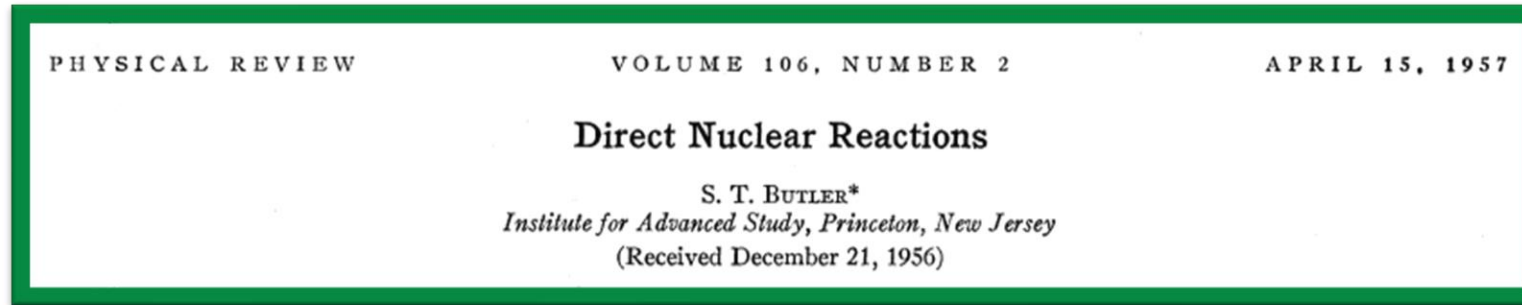


Single-particle spectroscopy of *exotic nuclei*:

The cases of ^{25}F and ^{207}Hg

Tsz Leung (Ryan) Tang
Argonne National Laboratory

Single-particle spectroscopy is a traditional probe. Used for more than 60 years.



Brief introduction to *Single-particle spectroscopy*

Single-particle spectroscopy

- Direct reactions that **ADD/REMOVE** a *single* nucleon.

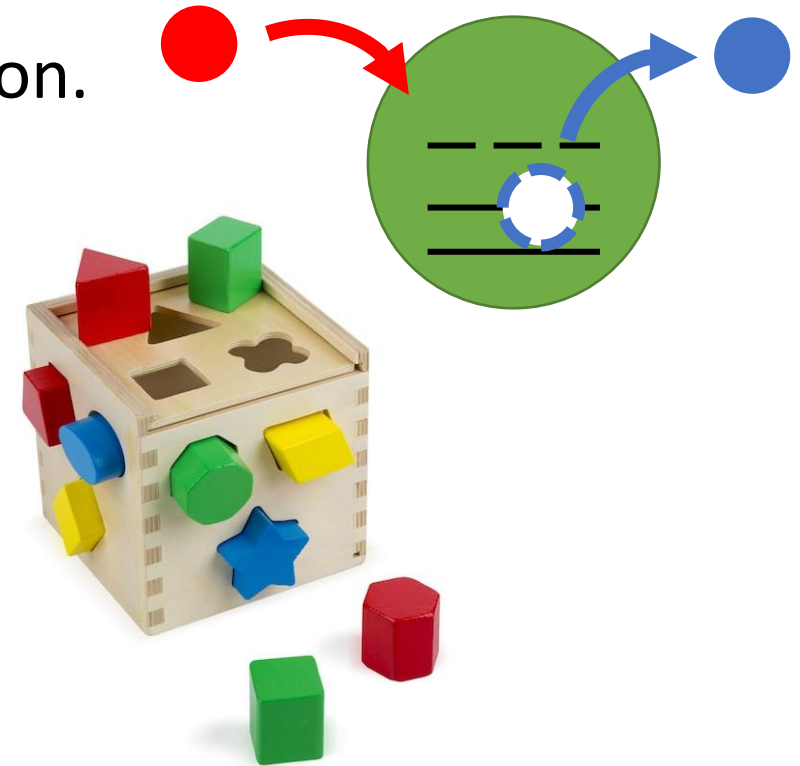
↳ single step process

- Measure **THREE** things

1. Energy level,
2. Angular distribution,
3. Yield / Cross section → Spectroscopic factor

- The power comes from the **SHELL MODEL**

- The total wave function is constructed by *single-particle wave function*



What is mean by single-particle?

The Hamiltonian of a nucleus A

$$H_A = \sum_i -\frac{\hbar^2}{2m_i} \nabla_i^2 + \sum_{i>j} V_{ij} + \sum_{i>j>k} V_{ijk} + \dots$$

Mean Field Approximation

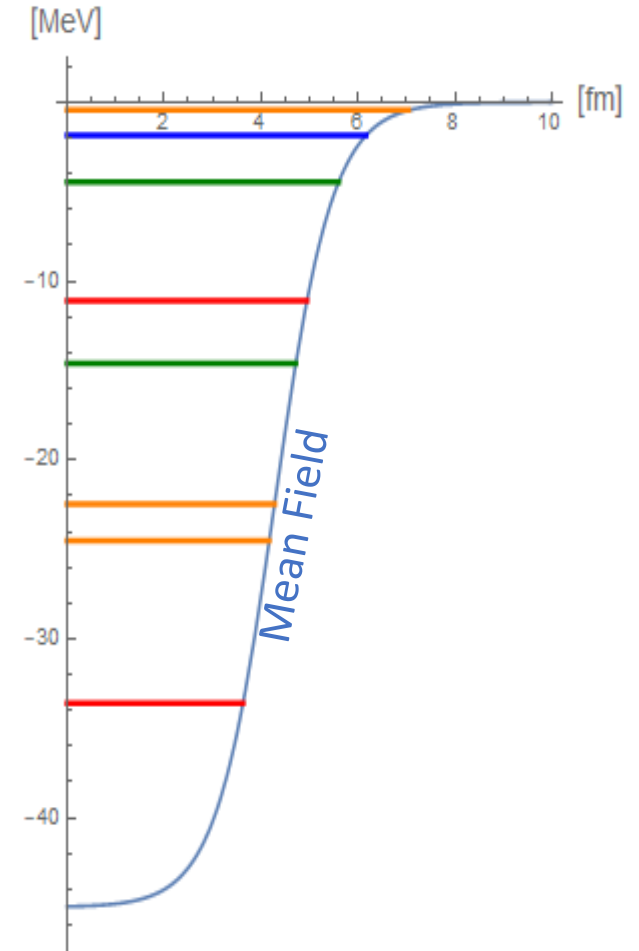
$$H_A = \left(\sum_i -\frac{\hbar^2}{2m_i} \nabla_i^2 + \mathbf{U} \right) + \left(\sum_{i>j} V_{ij} - \mathbf{U} \right) = \sum_i h_i + \sum_{i>j} R_{ij}$$

residual interaction

Single particle state

$$h_k \phi_{nlj} = \epsilon_{nlj} \phi_{nlj}$$

Single particle energy



Connection to single-particle spectroscopy

The Hamiltonian of a nucleus $A = B + k$

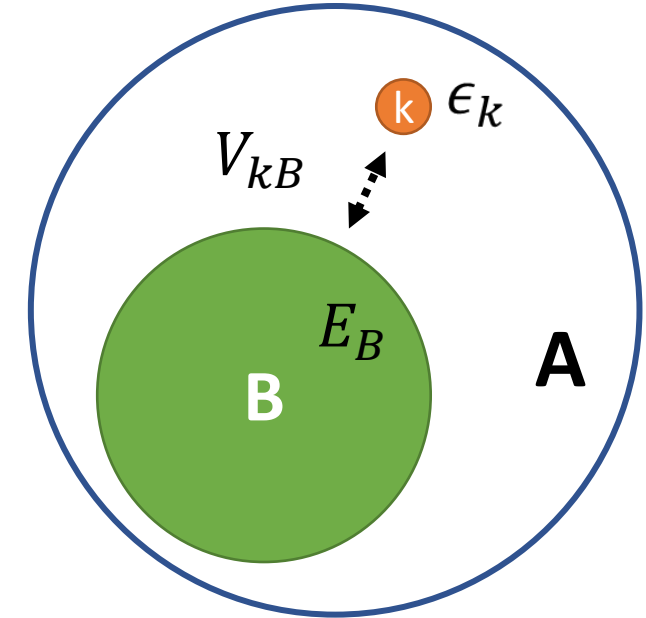
$$H_A = \sum_i h_i + \sum_{i \neq j} V_{ij} = h_k + H_B + V_{kB}$$

Total wave function of the nucleus

$$\Psi_{J_A} = \sum_{B'} \sum_{nlj} \beta_{nlj}(B', A) [\phi_{nlj} \Psi_{J_{B'}}]_{J_A}$$

$$h_k \phi_{nlj} = \epsilon_{nlj} \phi_{nlj}$$

$$H_B \Psi_{J_B} = E_{J_B} \Psi_{J_B}$$



Spectroscopic factor

$$S_{nlj}(B, A) = |\langle \Psi_B | a_{nlj}^+ | \Psi_A \rangle|^2 = |\beta_{nlj}(B, A)|^2$$

$$\langle \Psi_B | \Psi_A \rangle = \beta_{nlj}(B, A) \phi_{nlj}$$

Effective Single particle energy

$$\epsilon_{nlj} = \frac{\sum_B E_j |\beta_{nlj}(B, A)|^2}{\sum_B |\beta_{nlj}(B, A)|^2} \quad h_k \phi_{nlj} = \epsilon_{nlj} \phi_{nlj}$$

$$H_A \Psi_{J_A} = E_{J_A} \Psi_{J_A}$$

Intuitive picture via Single-particle spectroscopy

Spectroscopic factor

$$\langle \Psi_B | \Psi_A \rangle = \beta_{nlj}(B, A) \phi_{nlj}$$

- Similarity between 2 nuclear systems
- fraction of occupancy

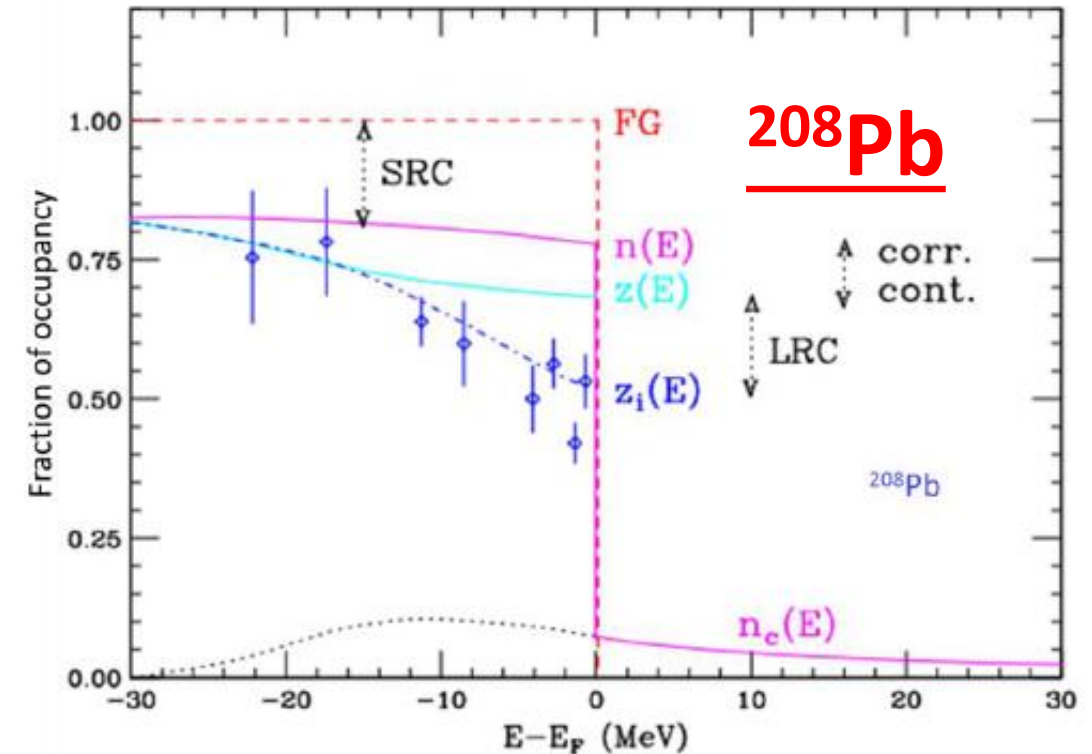
Effective single particle energy

$$\epsilon_{nlj} = \frac{\sum_B E_j |\beta_{nlj}(B, A)|^2}{\sum_B |\beta_{nlj}(B, A)|^2}$$

Weighted energy

- average out the perturbation by the residual interaction.
- monopole part of the Hamiltonian
- mean field picture.
- extract the **spin-orbit splitting** & more...

I. Sick, Quasi-free knockout reaction, QFS workshop at ECT, Trento, 2008.



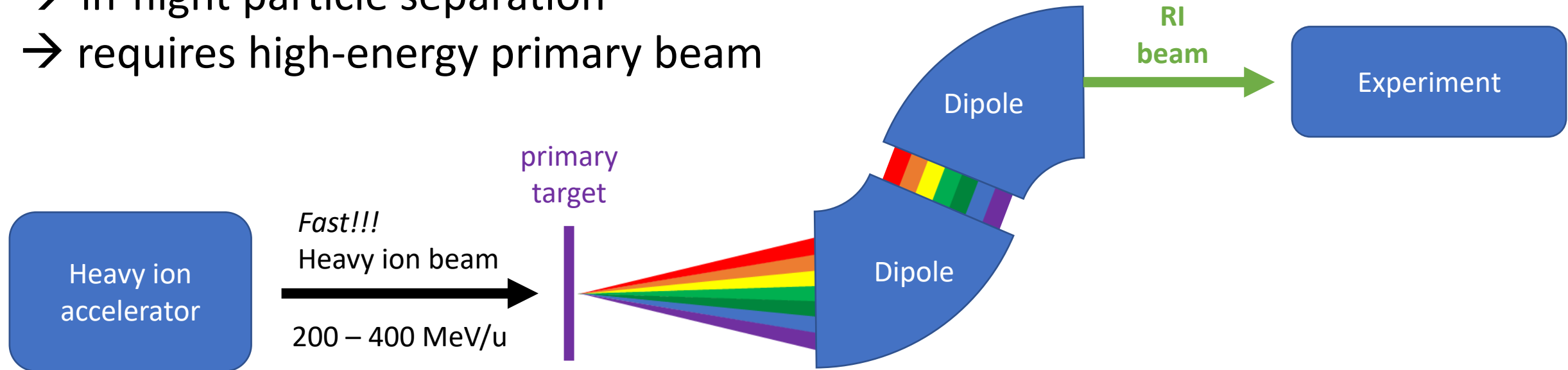
Fast beam case: $^{25}\text{F}(p, 2p)^{24}\text{O}$

Similarity between ^{25}F and $^{24}\text{O} + p$

Fast beam production

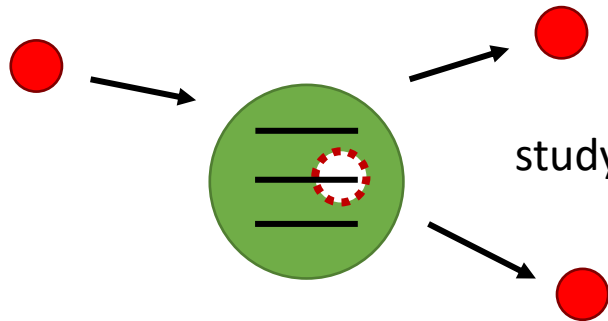
Projectile fragmentation method

- breaking up **heavy nucleus** into many **light nuclei**
- in-flight particle separation
- requires high-energy primary beam

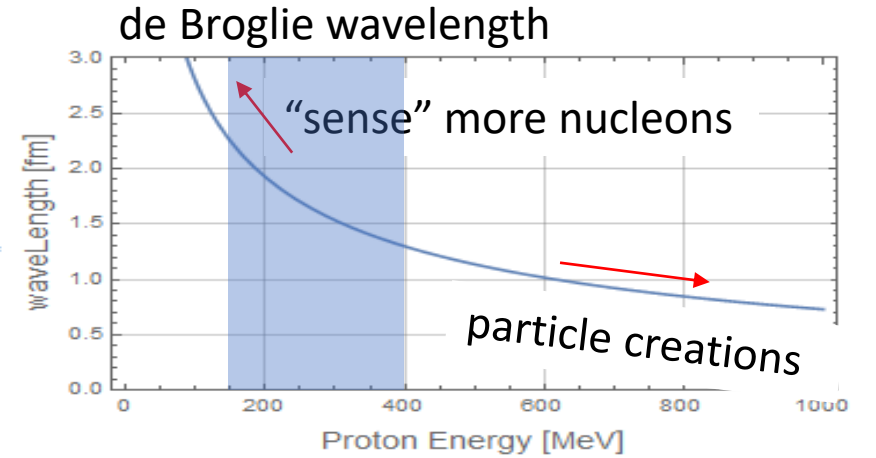


With fast beam, quasi-free ($p, 2p$) is a natural choice.

The $(p,2p)$ knockout Reaction

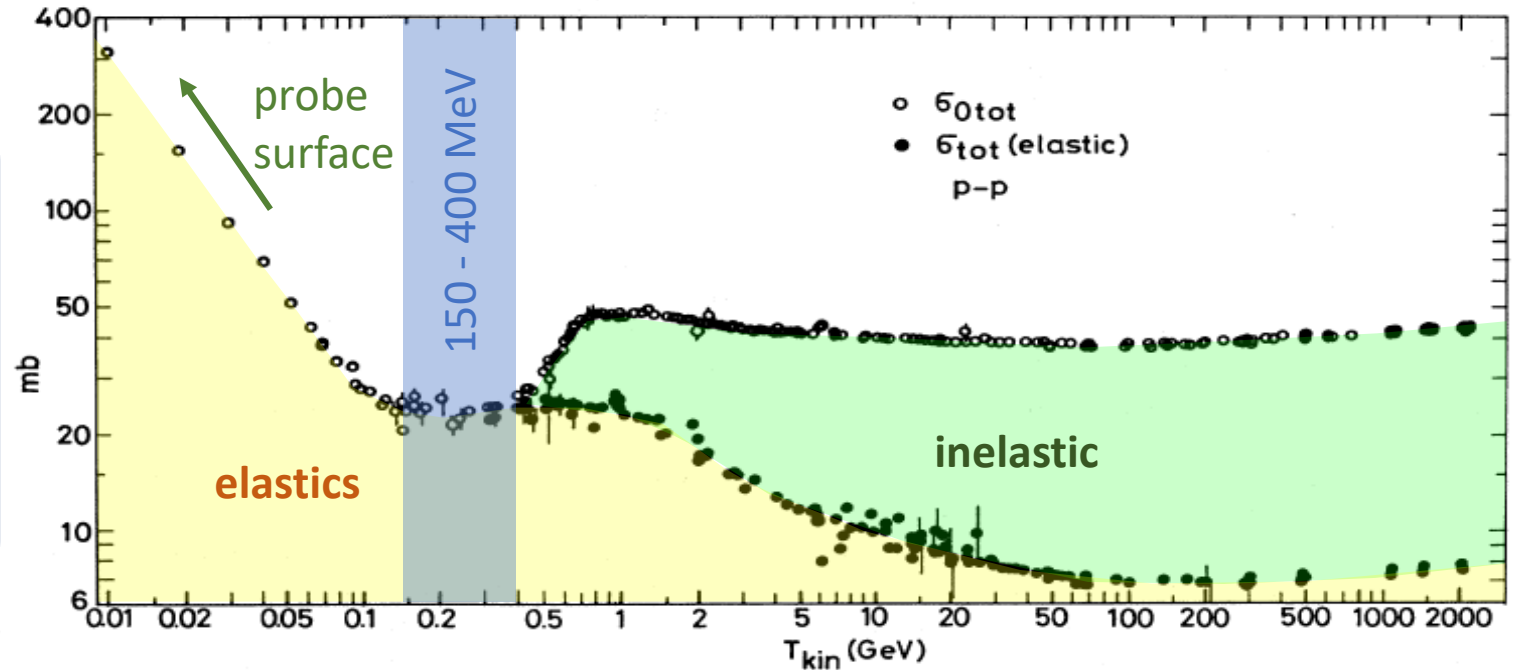


study the **fullness** of orbital



Ideal energy is 150 – 400 MeV

- Nucleus to be mostly transparent **Quasi-free!!**
- Nucleon motion is **frozen**
- Able to probe deeply bound states
- The reaction is **"CLEAN"**

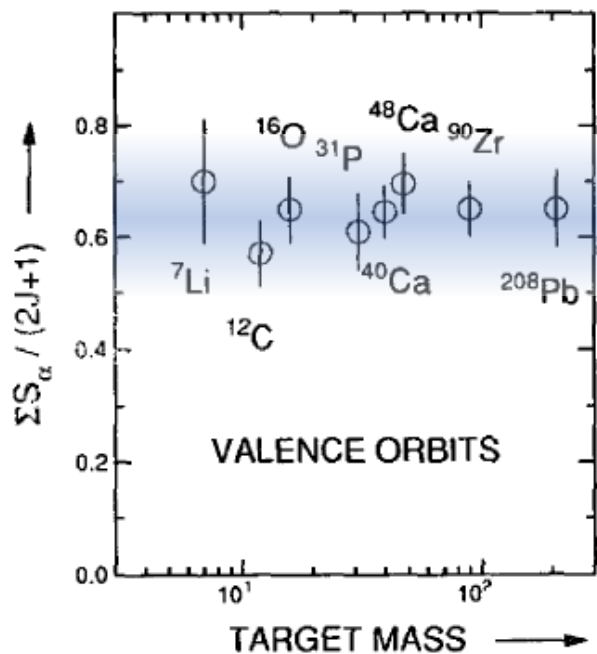


Rev. Mod. Phys. **65**, 47 (1993)

Quasi-free $(p,2p)$ is as “clean” as $(e,e'p)$

L. Lapikas, NPA **553**, 297 (1993)

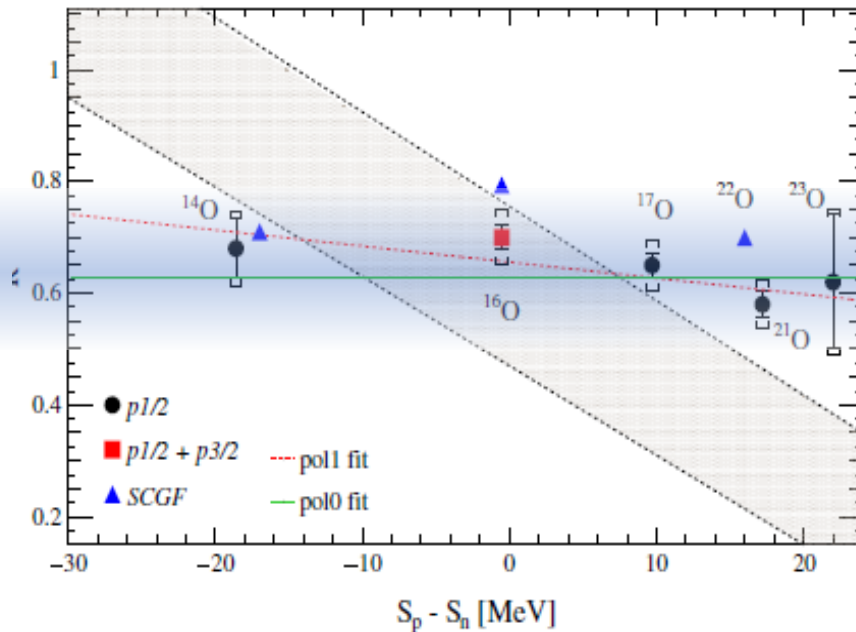
$(e,e'p)$



$(e,e'p)$: Coulomb interaction is well known.

L. Atar *et al.*, PRL **120**, 052501 (2018)

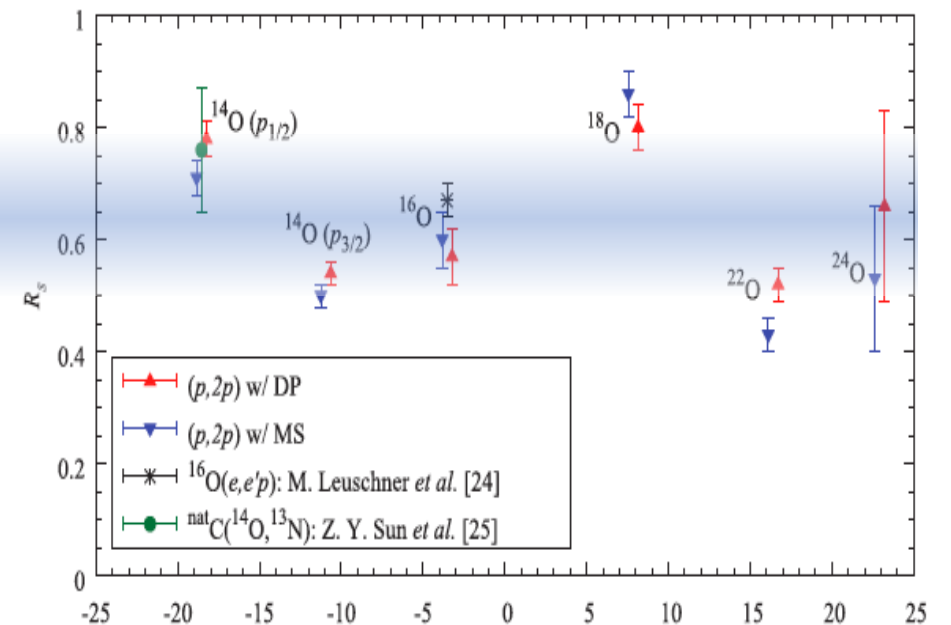
GSI, $(p,2p)$ @ 300 – 450 MeV/u



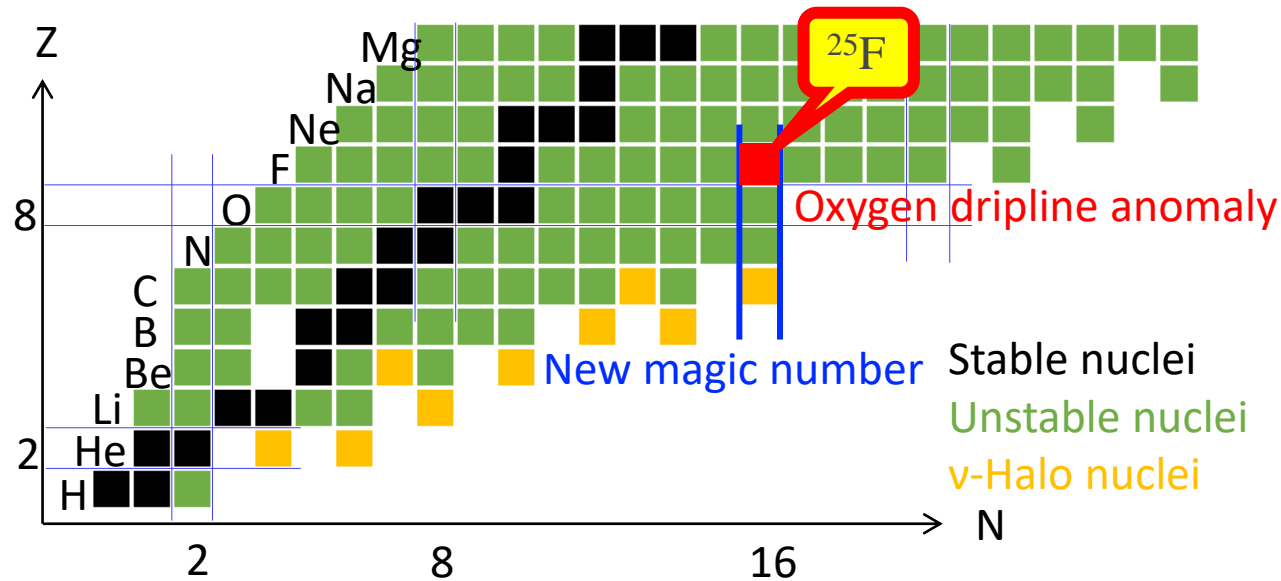
- Impulse approximation $\Delta S = S_p - S_n$ (MeV)
- delta interaction
- frozen nucleon motion

S. Kawase *et al.*, PTEP **021D01** (2018)

RCNP/RIKEN, $(p,2p)$ @ 200 - 250 MeV/u



$^{25}\text{F}(p,2p)^{24}\text{O}$: Oxygen dripline anomaly



Fluorine : 1-proton in sd shell.

→ support neutrons up to the p-f shell!

→ something interesting from that proton.

$^{25}\text{F}(p,2p)$

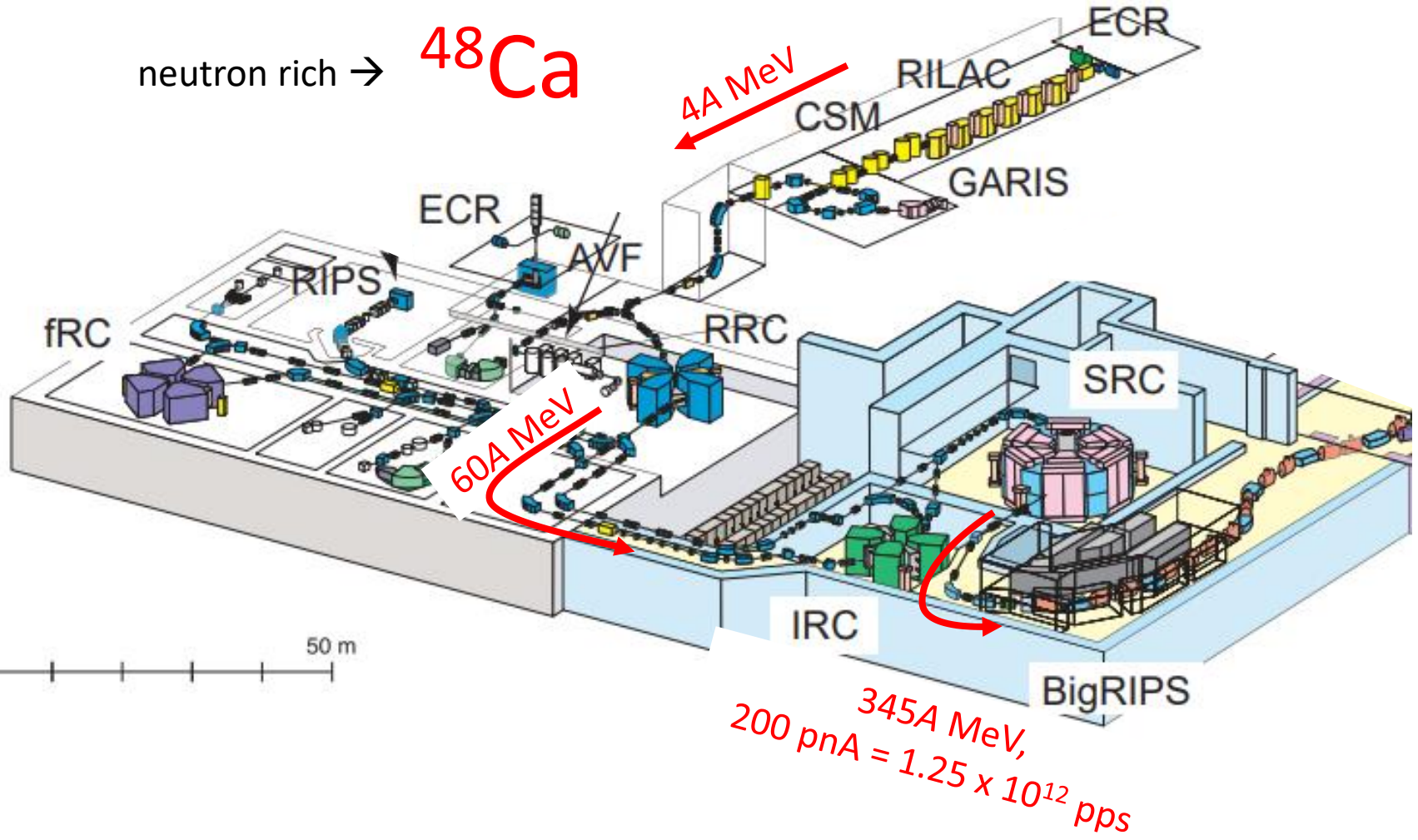
- ^{24}O is a doubly magic
→ the proton configuration mixing should be minimum.
- Knockout of the proton
→ study the neutron shell

$$\begin{aligned}
 S_{nlj}(O, F) &= |\langle \Psi_O | a_{nlj} | \Psi_F \rangle|^2 \\
 &= |\beta_{nlj}(O, F)|^2 \\
 &= |\langle \pi_O | a_{nlj} | \pi_F \rangle|^2 |\langle \nu_O | \nu_F \rangle|^2 \\
 &\quad \text{proton shell} \quad \text{neutron shell}
 \end{aligned}$$

^{25}F beam production (I)



Radioactive Ion Beam Factory (RIBF), RIKEN

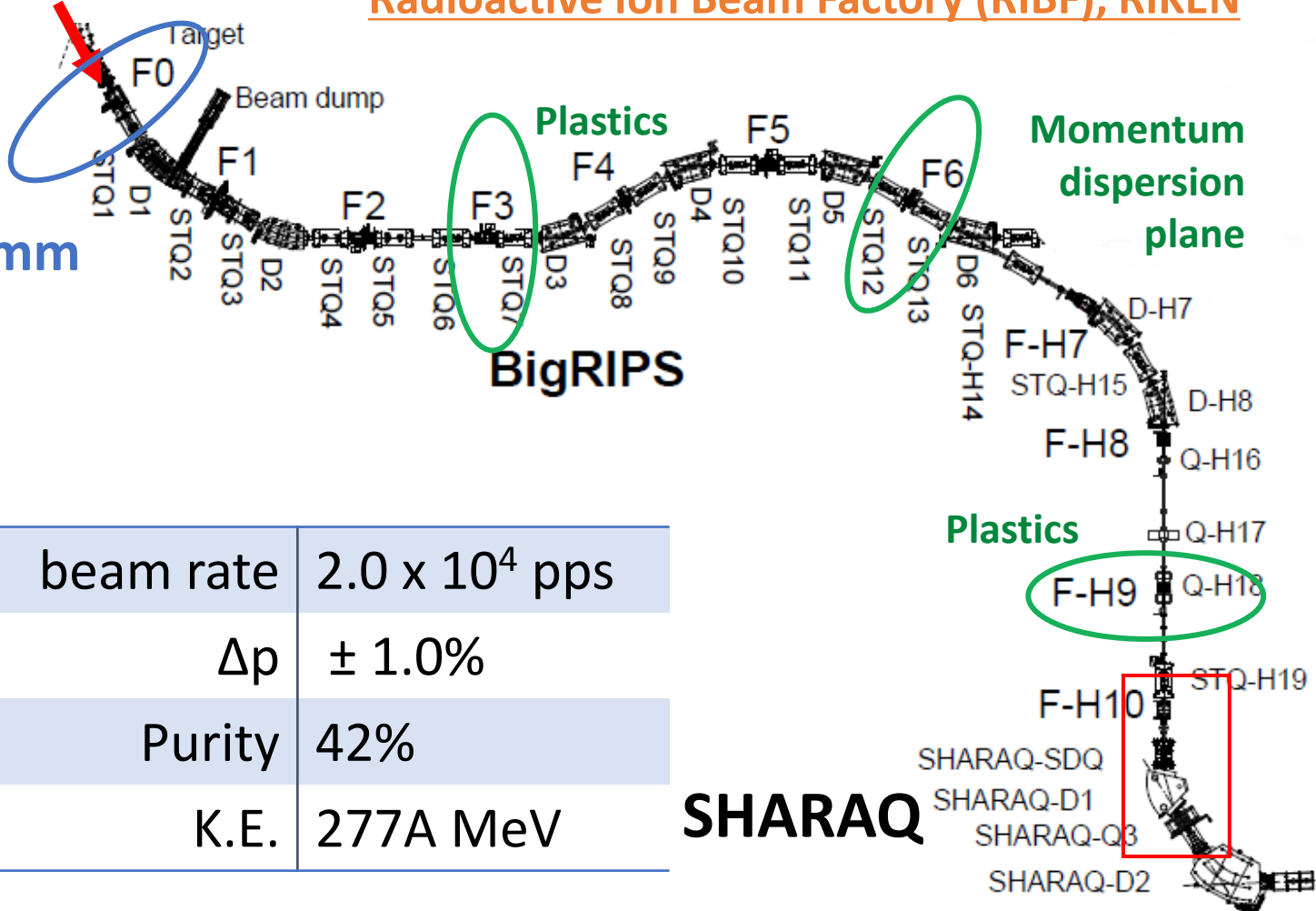


^{25}F beam production (II)

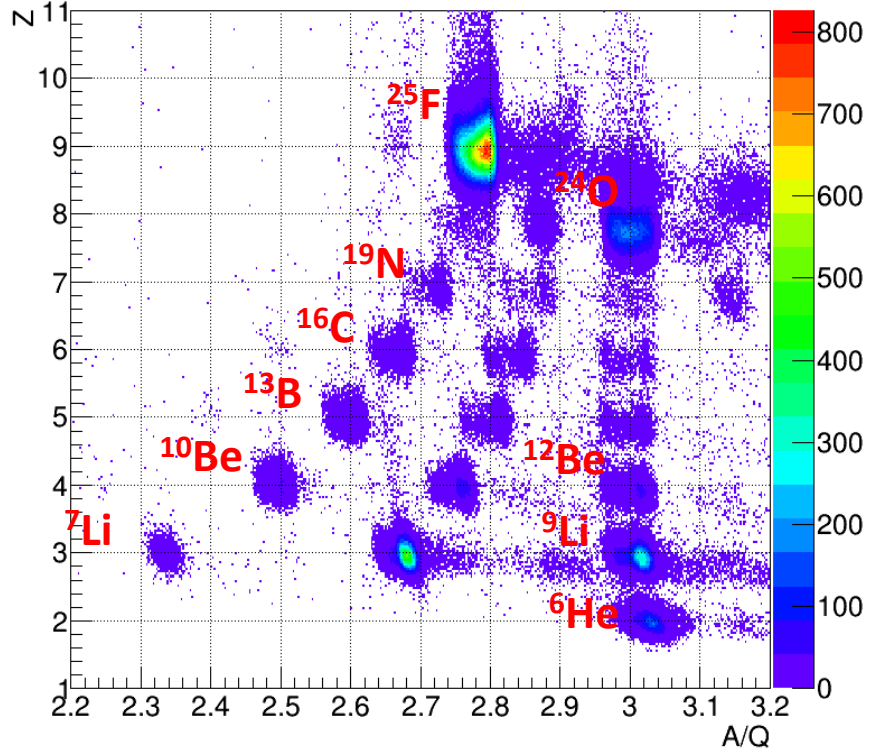
^{48}Ca
345A MeV
200pnA

^9Be 30mm

Radioactive Ion Beam Factory (RIBF), RIKEN



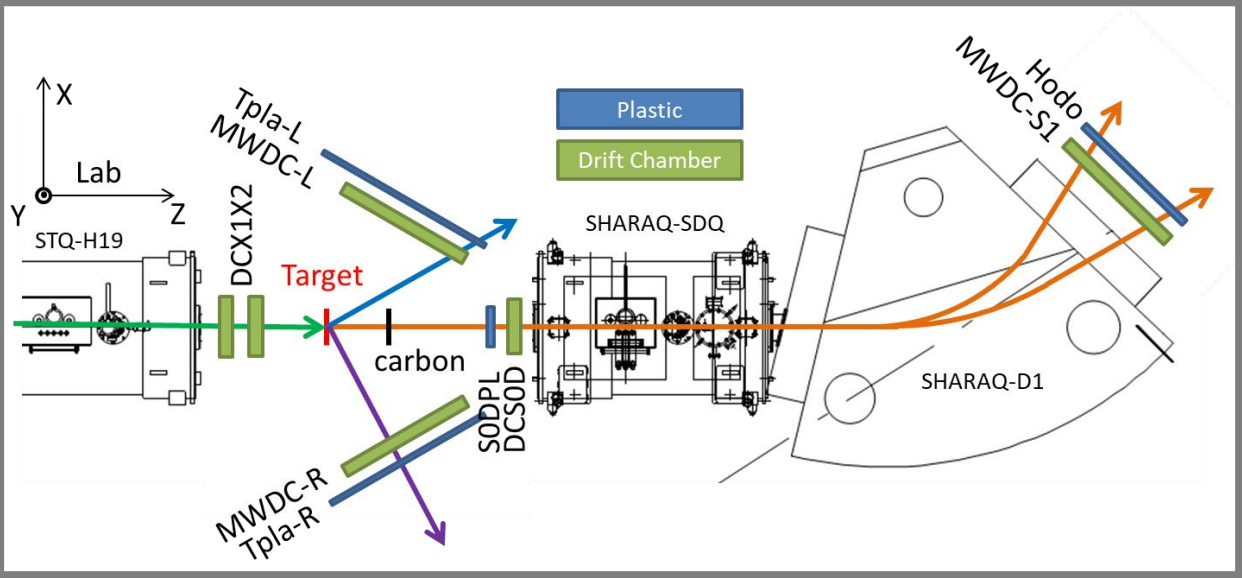
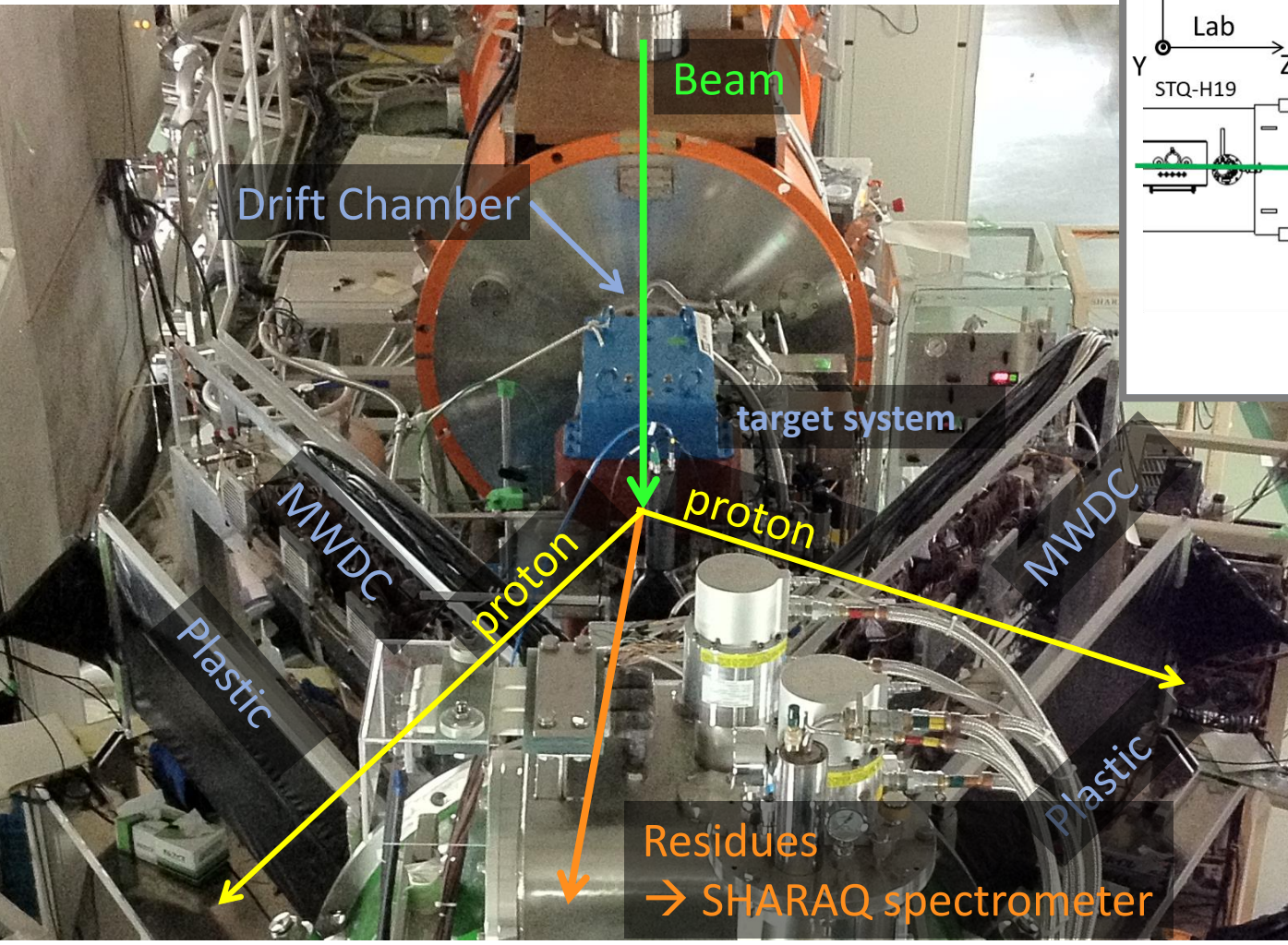
$\Delta E - \text{TOF} - \text{Bp}$ method



beam rate	2.0×10^4 pps
Δp	$\pm 1.0\%$
Purity	42%
K.E.	277A MeV

SHARQAQ

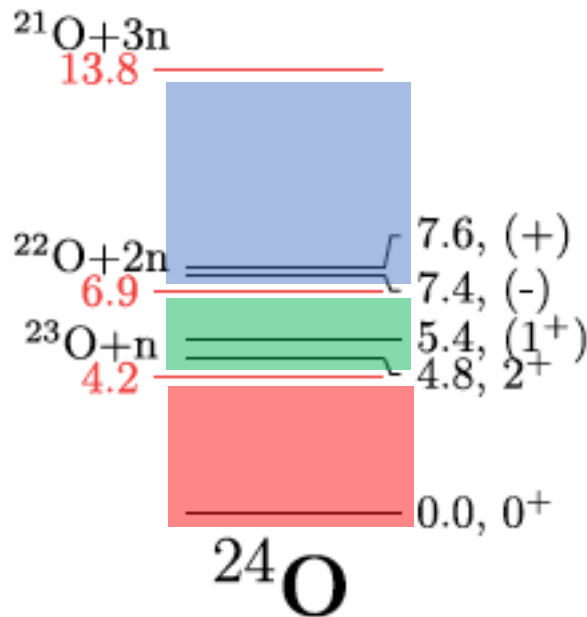
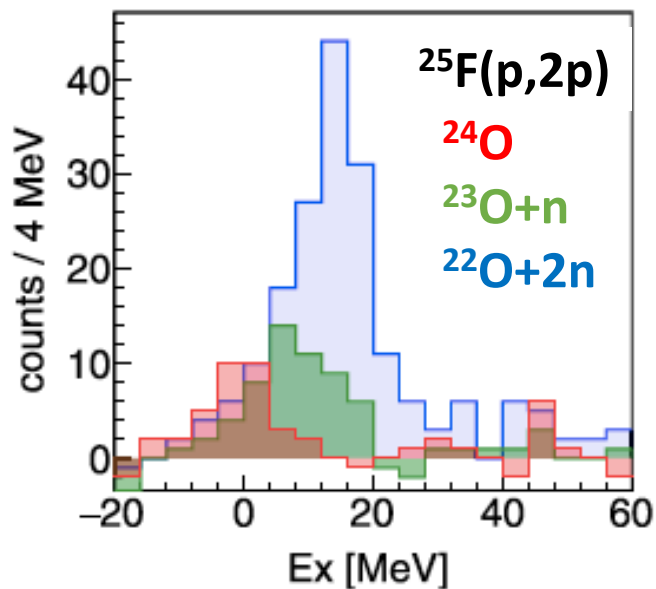
Experimental Setup



$$S_p = m_0 + m_2 - m_F$$

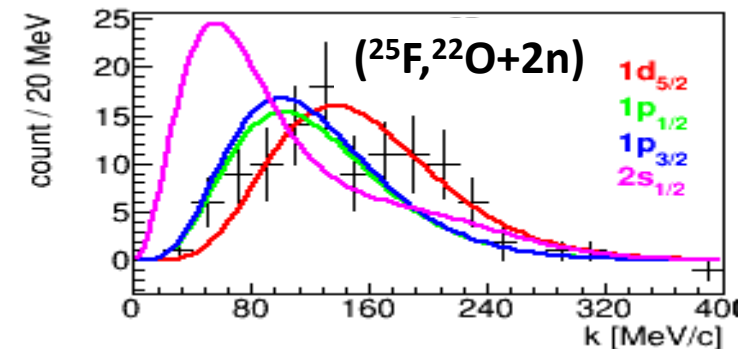
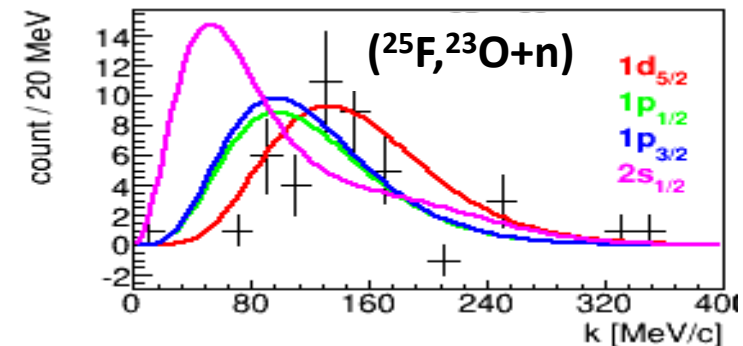
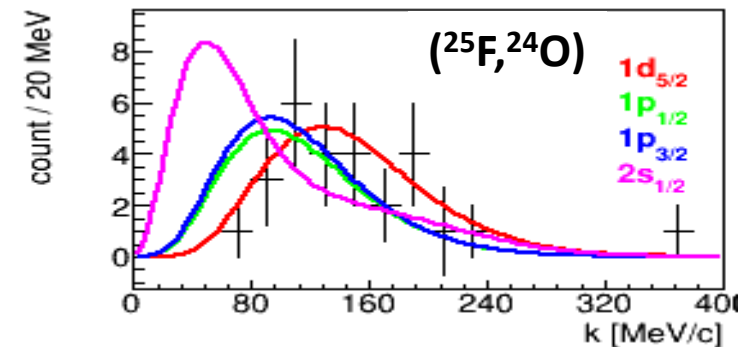
$$m_0 = m(\underbrace{\mathbb{P}_F + \mathbb{P}_p}_{\text{known}} - \underbrace{\mathbb{P}_1 - \mathbb{P}_2}_{\text{measure}})$$

Experimental Result



- $(^{25}\text{F},^{24}\text{O})$ is a single peak from $1d_{5/2}$ orbital.
- $(^{25}\text{F},^{23}\text{O}+n)$ is from sd-orbit \rightarrow no s-orbital \rightarrow $1d_{5/2}$ orbital
- Mean energy of $(^{25}\text{F},^{22}\text{O}+2n)$ is ~ 13 MeV
- shell gap between sd and p-shells = 12.7 MeV
 - p-orbital should dominate.

Momentum distribution



Spectroscopic factor and wave function of ^{25}F

$$\text{Spectroscopic factor} = \frac{\text{Exp. Cross Section}}{\text{DWIA Cross Section}}$$

(PIKOE, K. Ogata *et al.*, Osaka University, Japan)

	Residue	Orbital	SF
$^{25}\text{F}(p,2p)$	^{24}O	$1d_{5/2}$	0.36 ± 0.13
	$^{23}\text{O} + n$		0.65 ± 0.25

Optical potential :
Microscopic folding potential
(Melbourne G-matrix interaction)

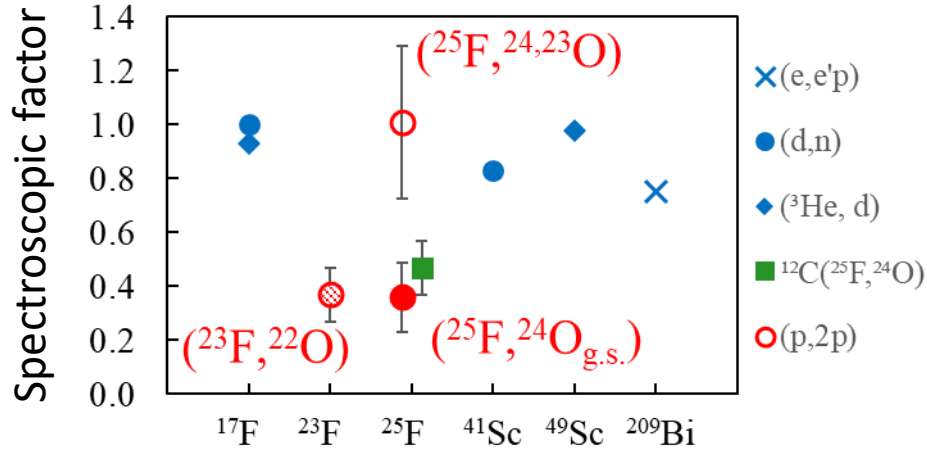
Bound state:
Woods-Saxon potential
 $r = 1.27A^{1/3}$ fm,
 $a = 0.67$ fm

Double magic

$$\text{Wave function of } ^{25}\text{F} = |^{25}\text{F}\rangle \approx |\pi 1d_{5/2}\rangle \otimes (\sqrt{0.36}|^{24}\text{O}_{\text{g.s.}}\rangle + \sqrt{0.65}|^{24}\text{O}^*\rangle + \dots)$$

Core of $^{25}\text{F} = \sim\mathbf{35\%}$ $^{24}\text{O}_{\text{g.s.}}$ and $\sim\mathbf{65\%}$ ^{24}O excited states.

Why the G.S. spectroscopic factor is so small?



- In double magic + p nuclei, the quenching is small.
- The proton in ^{25}F is almost stay in $d_{5/2}$ shell.
 → it seems that **proton is in SPS**
 - There are 65% excited ^{24}O in ^{25}F core.
 → something on **the neutron side.**

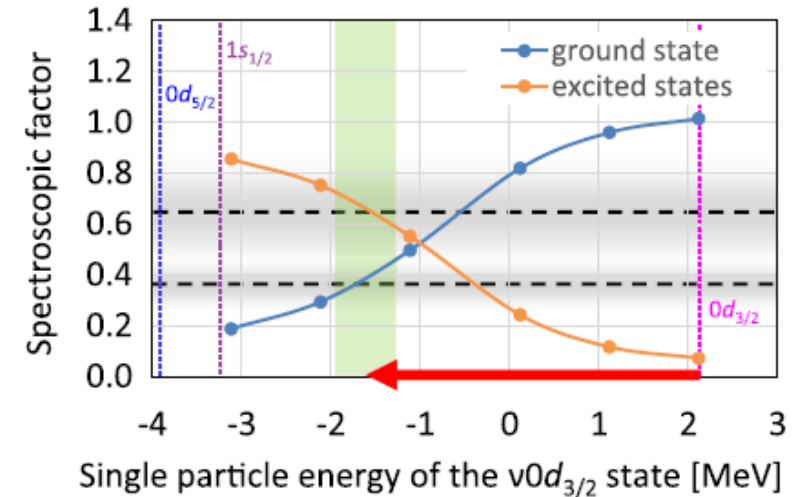
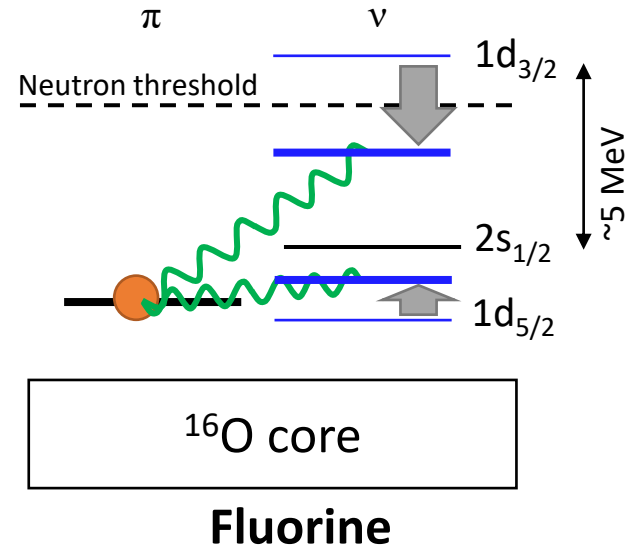
Possible Mechanism:

Type-I shell evolution driven by **tensor force**

T. Otuska *et al.*, J. Phys. G: Nucl. Part. Phys. **43** (2016) 024009

Consequence are:

- Increase neutron configuration mixing
- Disappear of $N = 16$ magicity.
- Long Fluorine neutron dripline



Can shell model calculations explain the result?

	Residue	Orbit	SF _{exp}	SF(SFO)	SF(USDB)	SF(SDPF-MU)
²⁵ F(p,2p)	²⁴ O	1d _{5/2}	0.36 ± 0.13	0.9	1.01	0.95
	²³ O + n		0.65 ± 0.25	0.1	-	-
		Model space		p-sd	sd	sd-pf
		Reference		T. Suzuki <i>et al.</i> , PRC 67 (2003) 044302	B. A. Brown <i>et al.</i> , PRC 74 (2006) 034315	Y. Utsno <i>et al.</i> , RPC 86 (2012) 051301

with 0, 1, 2 ħω

- Give almost unity of ground state
- Produce no/little fragmentation.

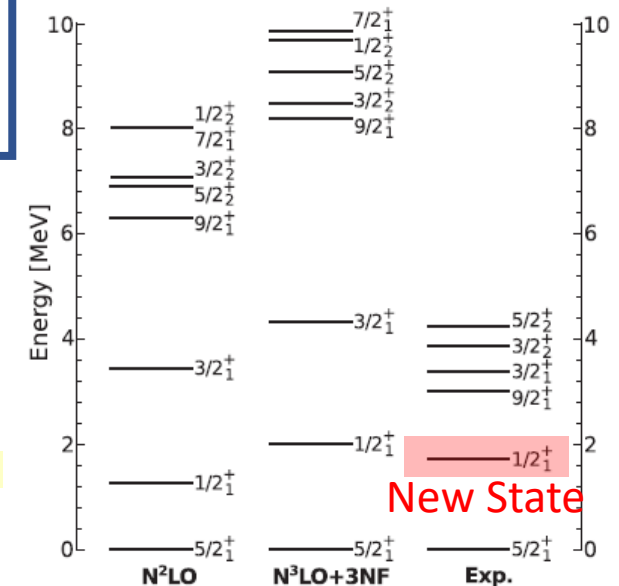
May be...

- The tensor force is not strong enough?
- Real 3N force is needed?
- Strong deformation of ²⁵F?

TABLE I. Partial norms of wave functions using up to 3p-2h amplitudes on top of the ²⁴O reference state. The interaction is the newly optimized chiral interaction at third order (N²LO) [31].

	1p	2p-1h	3p-2h
5/2 ₁ ⁺	0.63	0.30	0.07
1/2 ₁ ⁺	0.56	0.36	0.08
9/2 ₁ ⁺	0.00	0.74	0.26

in-beam γ-ray spectroscopy



Zs. Vajta *et al.*, PRC **89**, 054323 (2014)

PHYSICAL REVIEW LETTERS **124**, 212502 (2020)**How Different is the Core of ^{25}F from $^{24}\text{O}_{\text{g.s.}}$?**

T. L. Tang^{1,2,*} T. Uesaka² S. Kawase^{1,†} D. Beaumel,³ M. Dozono,² T. Fujii,¹ N. Fukuda,² T. Fukunaga,⁴
 A. Galindo-Uribarri,⁵ S. H. Hwang,^{6,‡} N. Inabe,² D. Kameda,² T. Kawahara,⁷ W. Kim,⁶ K. Kisamori,¹ M. Kobayashi,¹
 T. Kubo,² Y. Kubota,^{1,§} K. Kusaka,² C. S. Lee,¹ Y. Maeda,⁸ H. Matsubara^{2,||} S. Michimasa,¹ H. Miya,¹ T. Noro,⁴
 A. Obertelli,^{2,9,§} K. Ogata^{10,11} S. Ota,¹ E. Padilla-Rodal¹² S. Sakaguchi⁴ H. Sakai,² M. Sasano,² S. Shimoura¹ S.
 S. Stepanyan,⁶ H. Suzuki,² M. Takaki,¹ H. Takeda,² H. Tokieda,¹ T. Wakasa,⁴ T. Wakui,^{13,¶} K. Yako,¹ Y. Yanagisawa,²
 J. Yasuda,⁴ R. Yokoyama,¹ K. Yoshida,² K. Yoshida,^{10,**} and J. Zenihiro²



Slow beam case: $^{206}\text{Hg}(d,p)^{207}\text{Hg}$

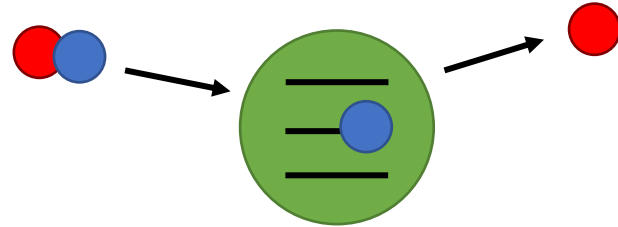
Transfer reaction

Solenoidal spectrometer

Transfer reaction

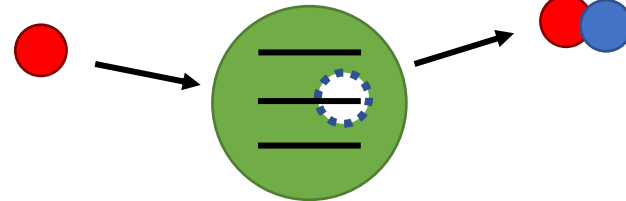
(d,p) neutron *removal*

study the **emptiness** of orbital

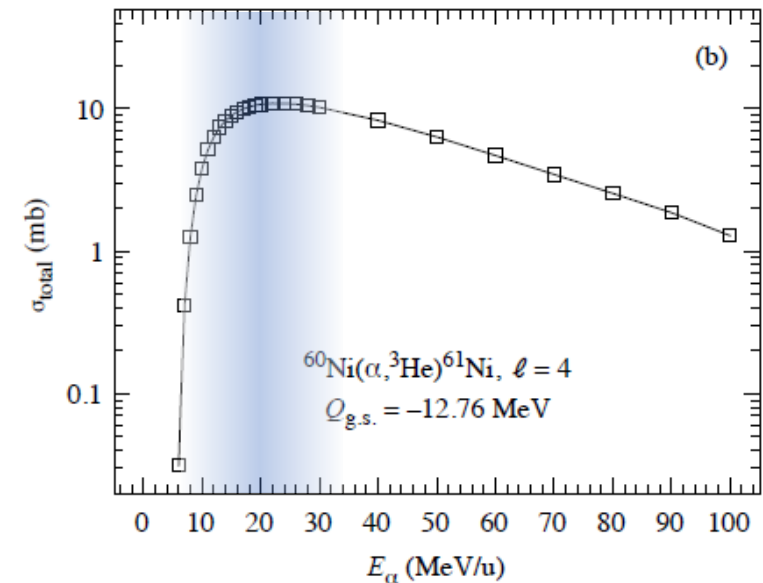
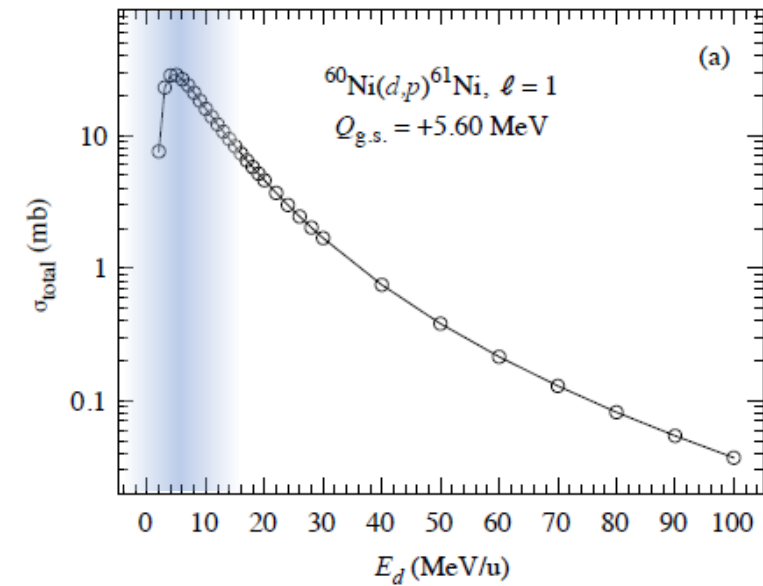


(p,d) neutron *adding*

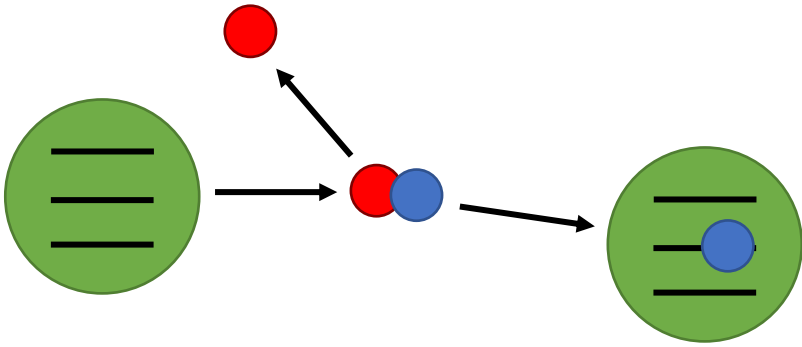
study the **fullness** of orbital



- Momentum matching
 - **Slow beam**
 - 5 ~ 20 MeV/u for maximum cross section
- Reaction on surface.
- Only 2 degree of freedoms: θ_{cm} and E_x

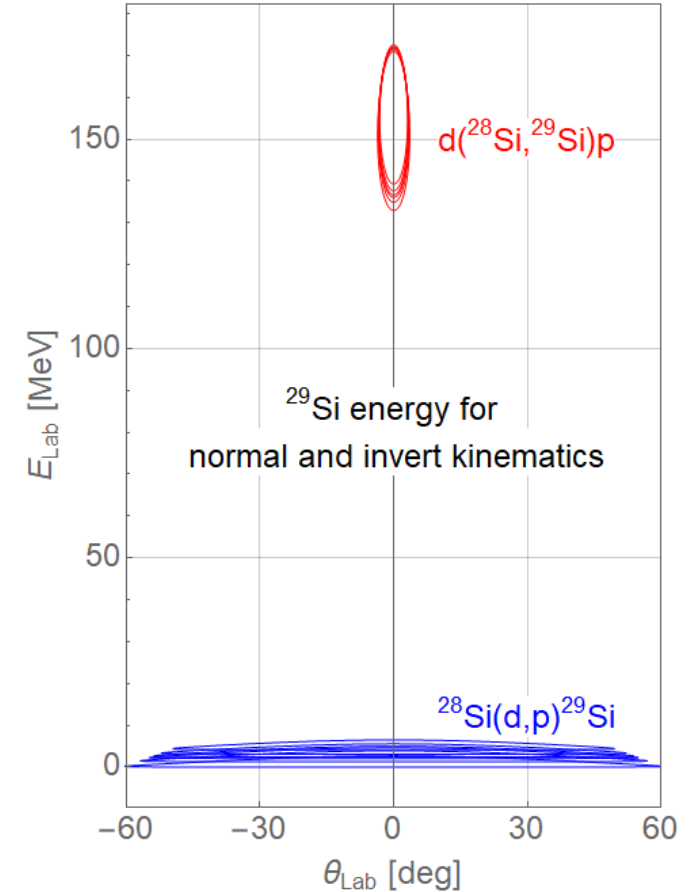
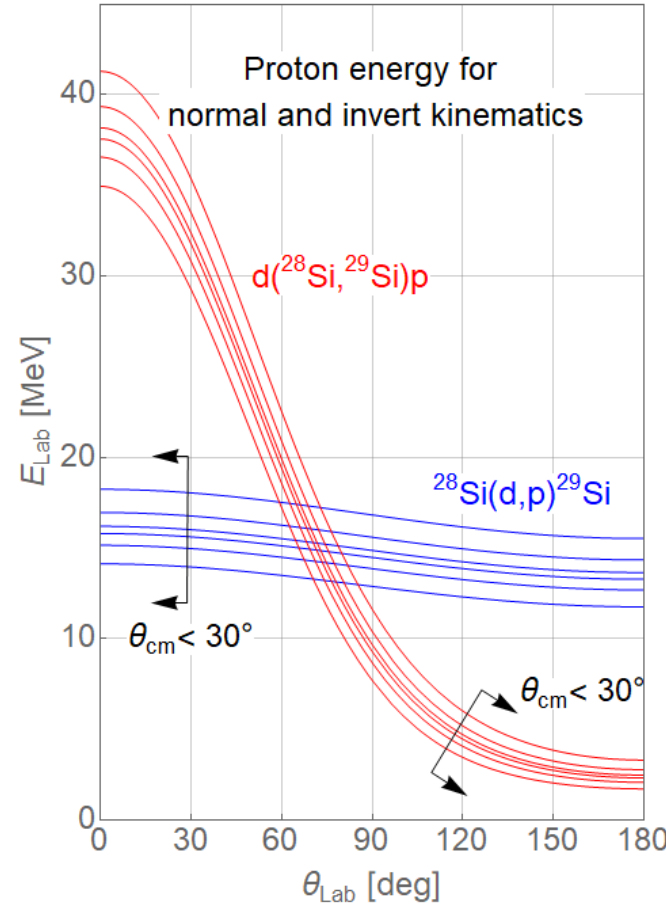


Inverse Kinematics for transfer reaction

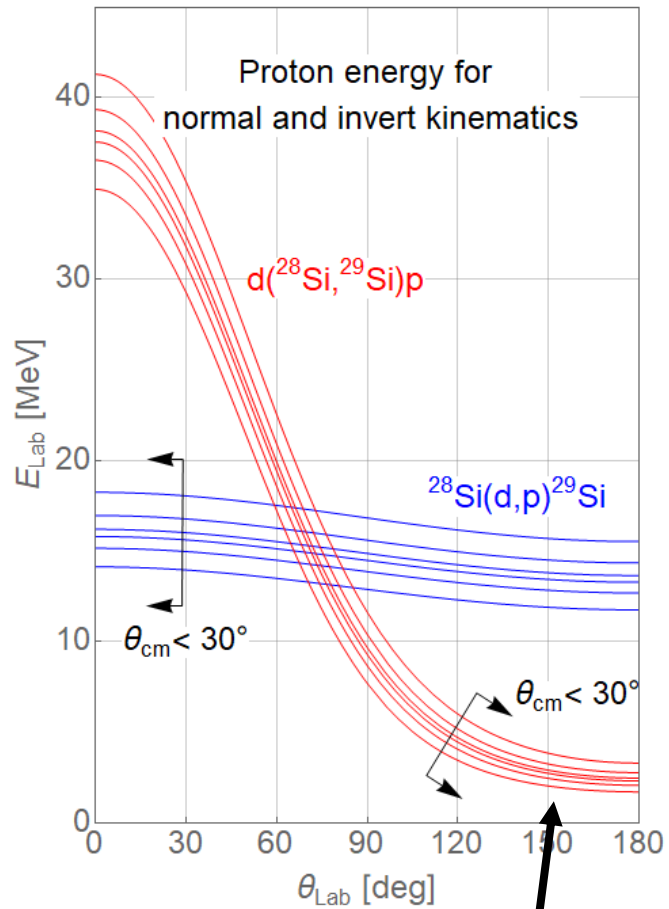


Advantage of inverse kinematics:

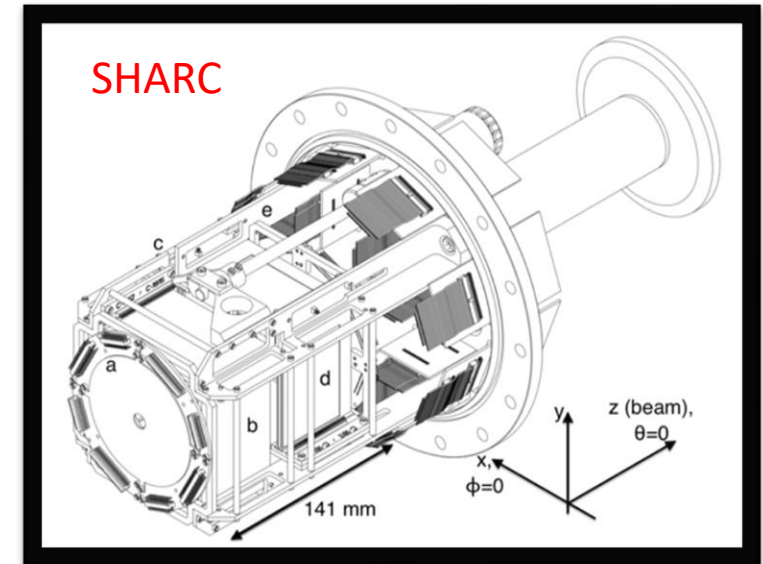
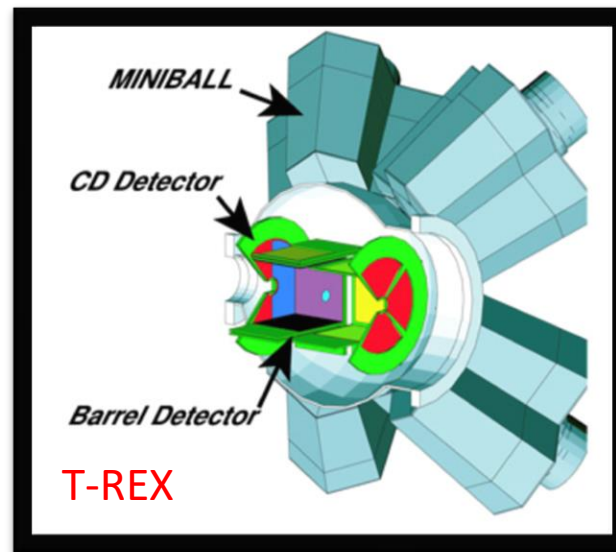
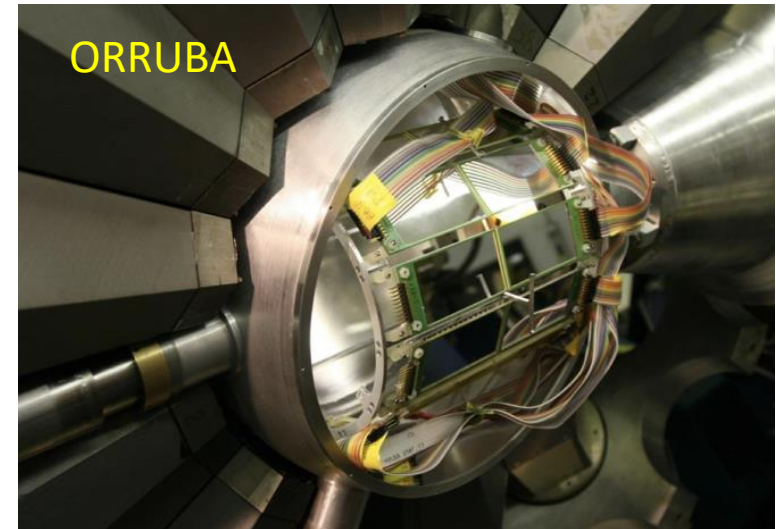
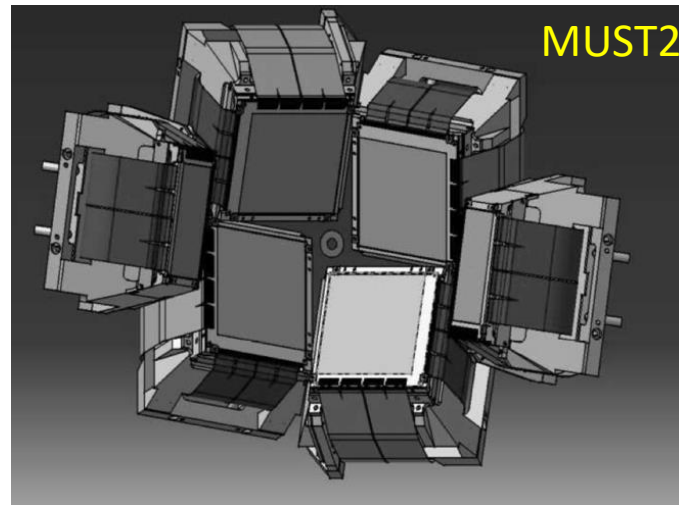
- Detection of the residual nucleus
- Smaller KE for the light particle
- Larger spread of θ_{Lab} for $\theta_{cm} \in (0^\circ, \sim 30^\circ)$



Some detectors systems

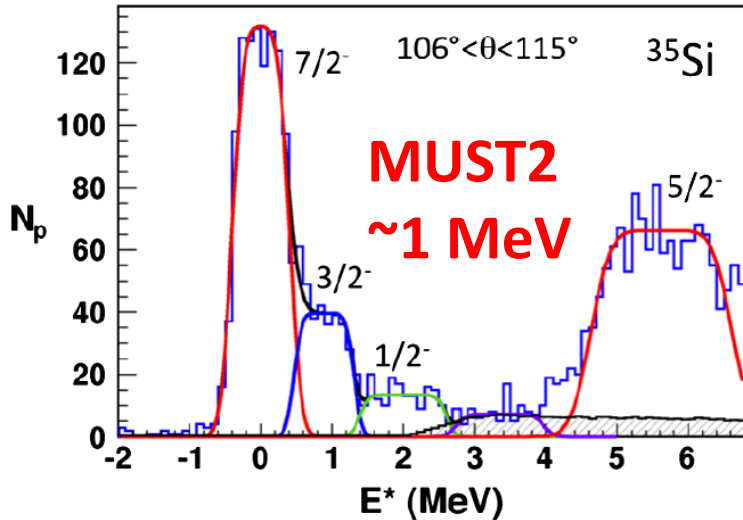


Cross section concentrates on $\theta_{\text{cm}} < 30^\circ$

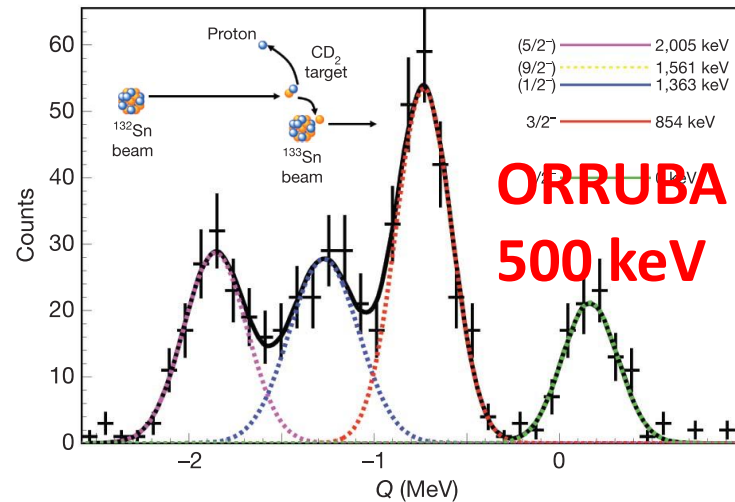


Energy resolution is challenging...

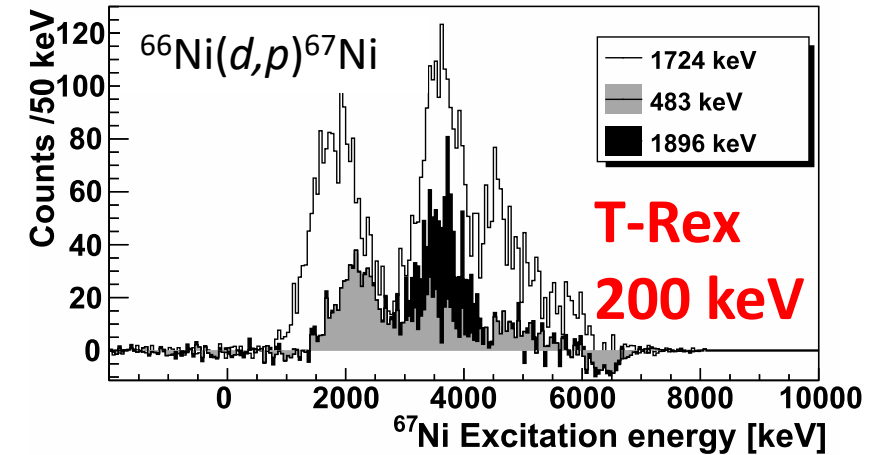
G. Burgunder et al., PRL **112**, 042502 (2014)



K. L. Jones et al., Nature **465**, 454 (2010)

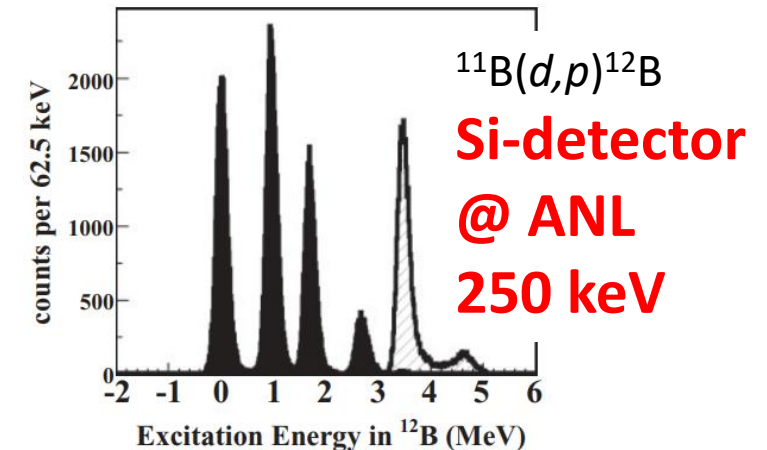


J. Diriken et al., Phys. Rev. C **91**, 054321 (2015)



- Resolution typically around $\sim 200 - 500$ keV FWHM
- They are near the target
 \rightarrow angular resolution is limited.

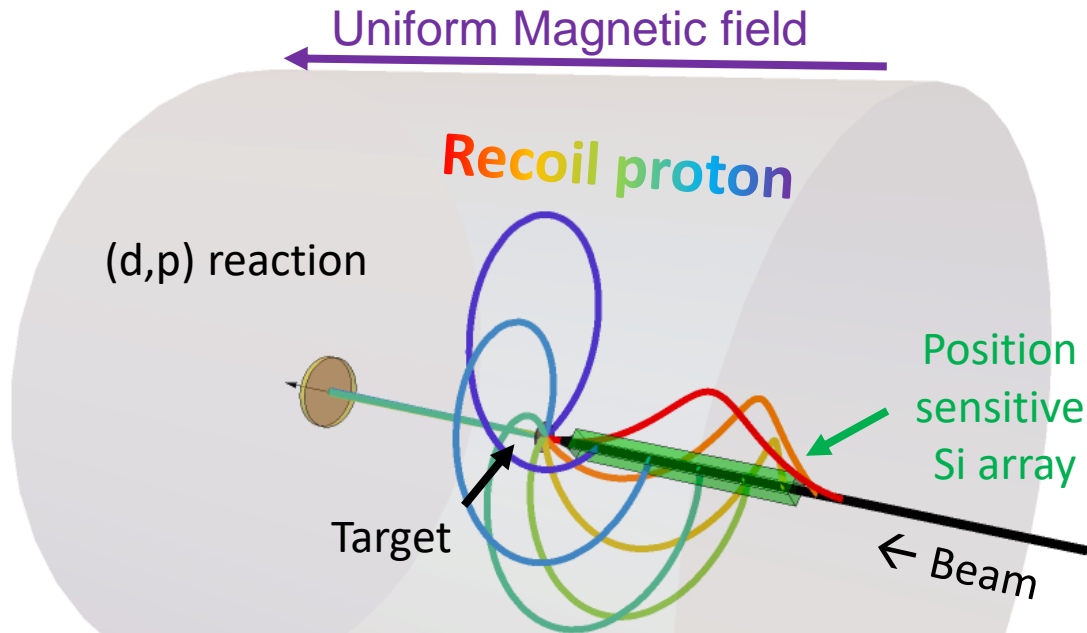
It is better to have resolution ~ 100 keV FWHM.



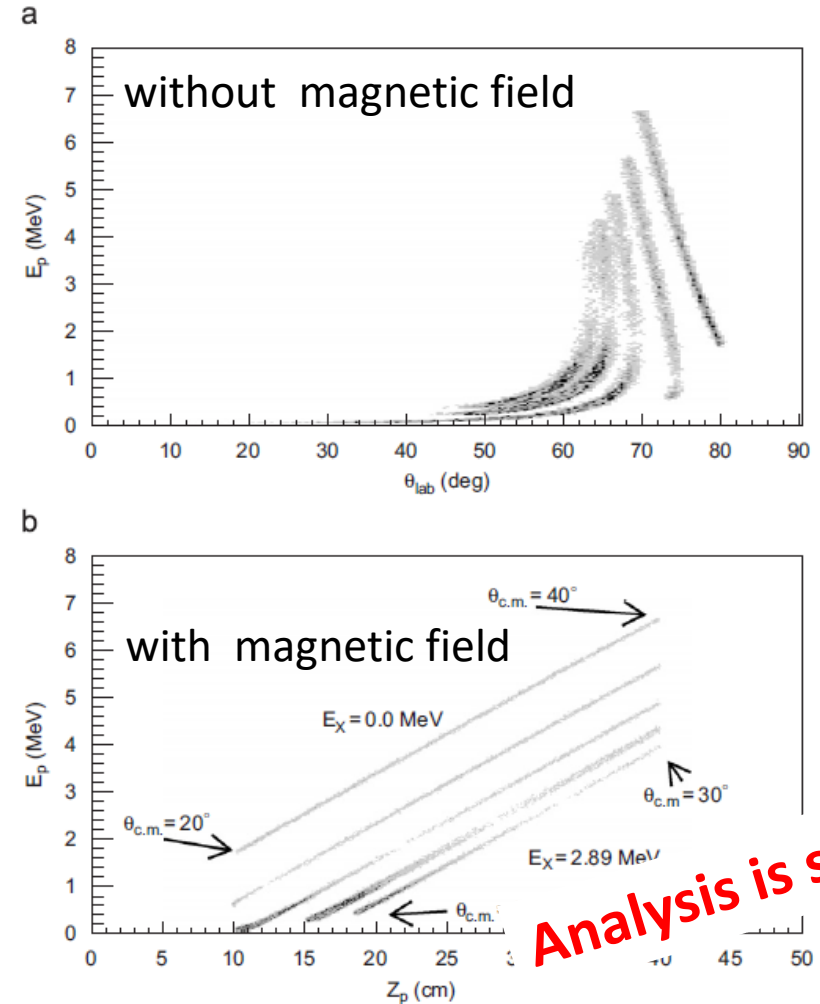
H. Y. Lee et al., PRC **81**, 015802 (2010)

The Idea of Solenoidal Spectrometer

A.H. Wuosmaa, J.P. Schiffer, B.B. Back, C.J. Lister, K.E. Rehm, NIM A **580** (2007) 1290



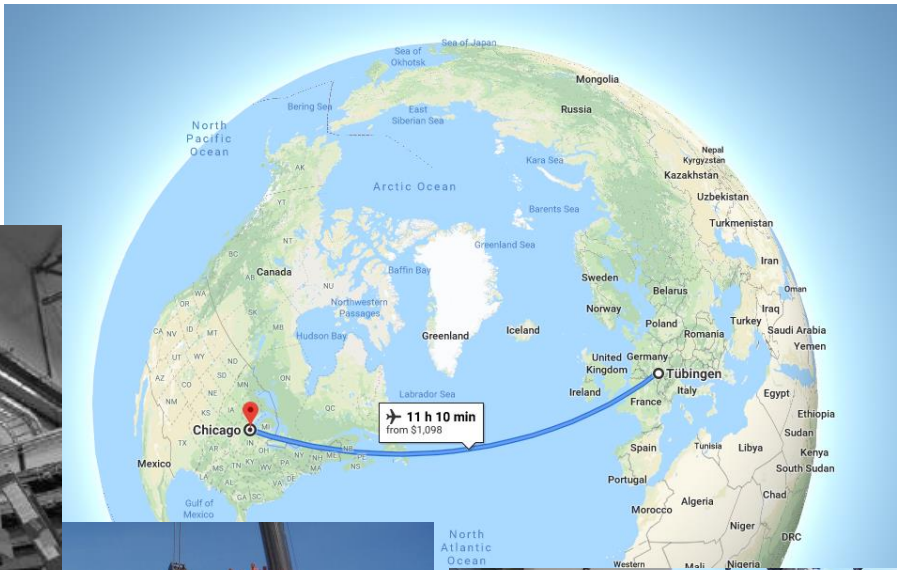
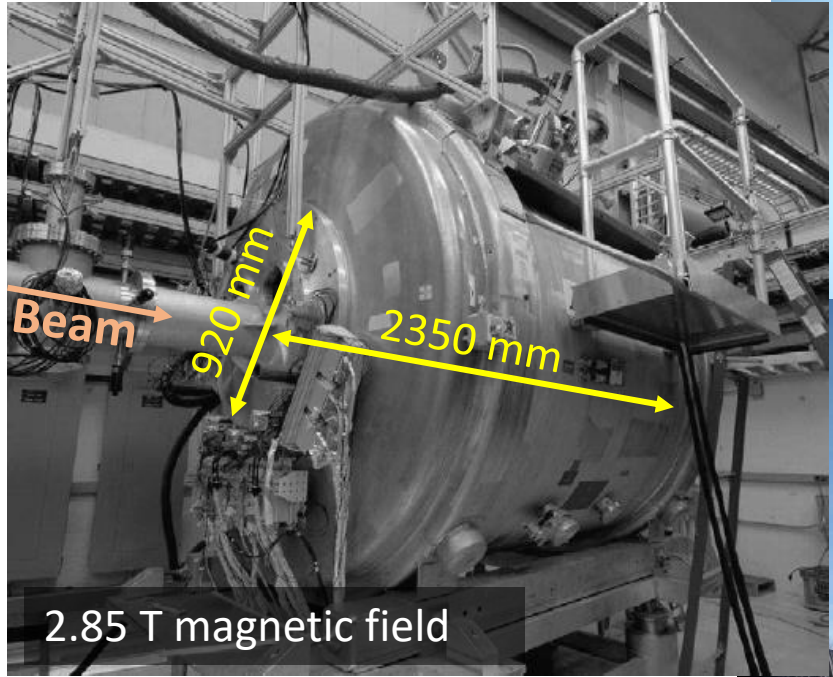
- Large Acceptance
- Good angular resolution
- Energy resolution ~ 100 keV FWHM



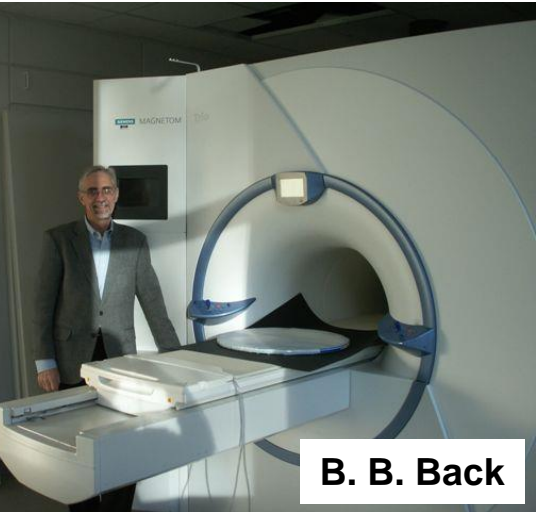
Analysis is simple!!

HELIOS in ANL

Decommissioned
Magnetic Resonance Imaging device

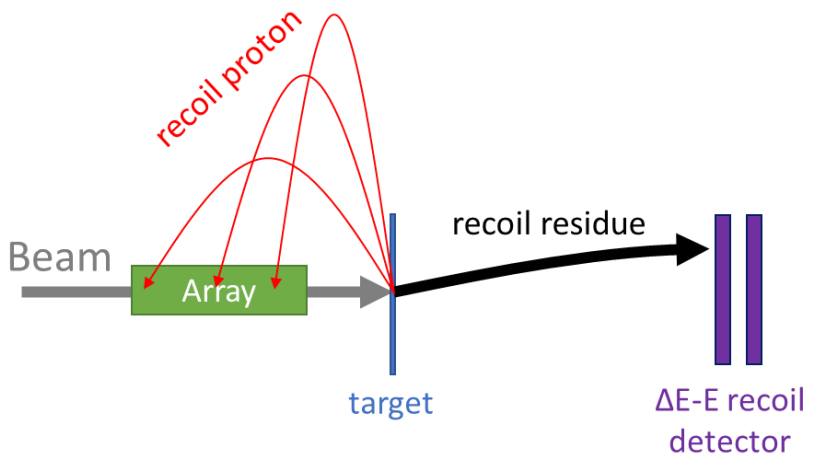
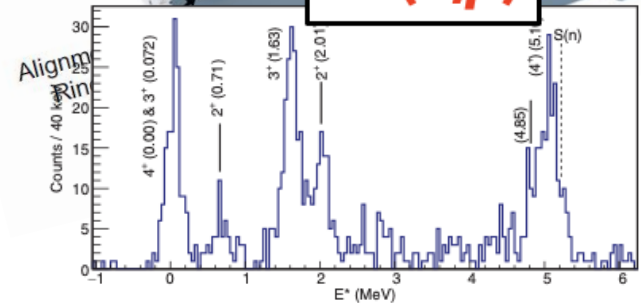
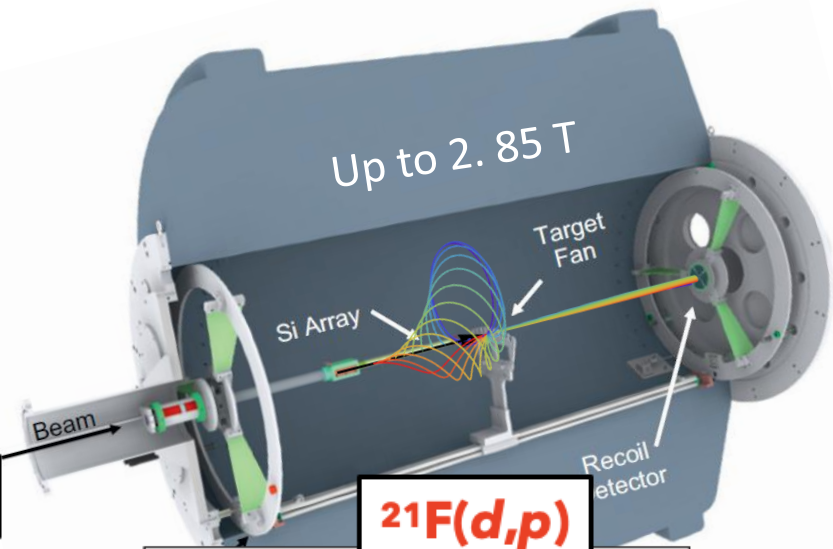
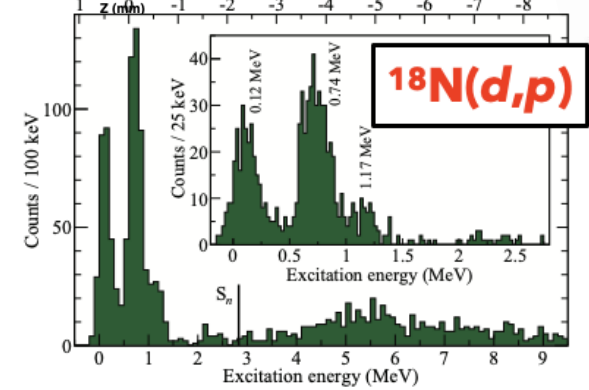
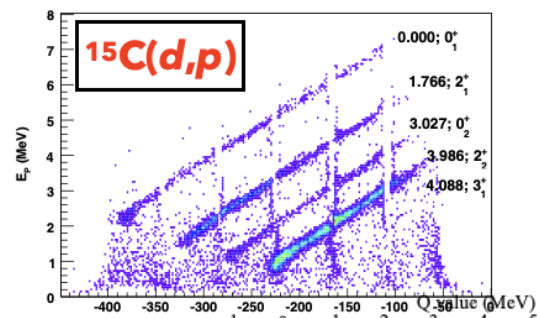


Tübingen, Monday, Nov 6, 2006

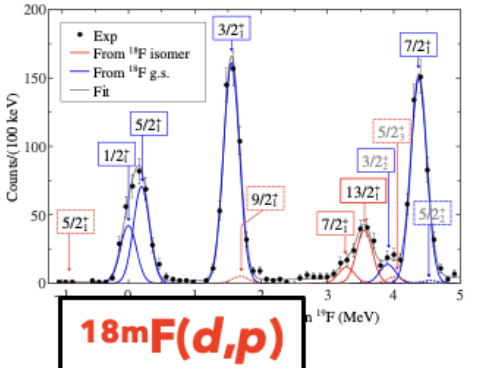
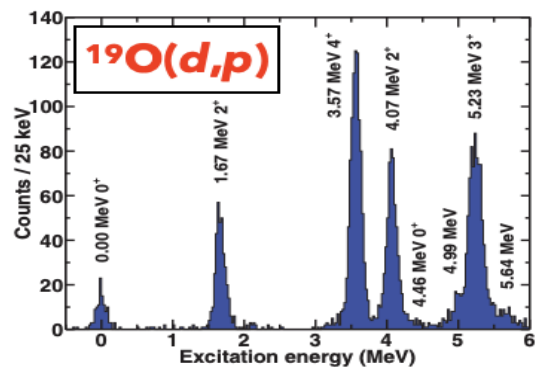
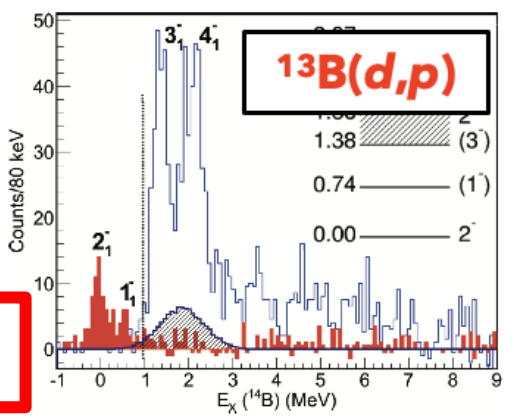


The (d,p) reactions with HELIOS

						16Ne	17Ne	18Ne	19Ne	20Ne	21Ne	22Ne	23Ne	24Ne	25Ne		
						14F	15F	16F	17F	18F	19F	20F	21F	22F	23F	24F	
						12O	13O	14O	15O	16O	17O	18O	19O	20O	21O	22O	23O
10N	11N	12N	13N	14N	15N	16N	17N	18N	19N	20N	21N	22N					
9C	10C	11C	12C	13C	14C	15C	16C	17C	18C	19C	20C	21C					
8B	9B	10B	11B	12B	13B	14B	15B	16B	17B	18B	19B	20B					
7Be	8Be	9Be	10Be	11Be	12Be	13Be	14Be	15Be	16Be								

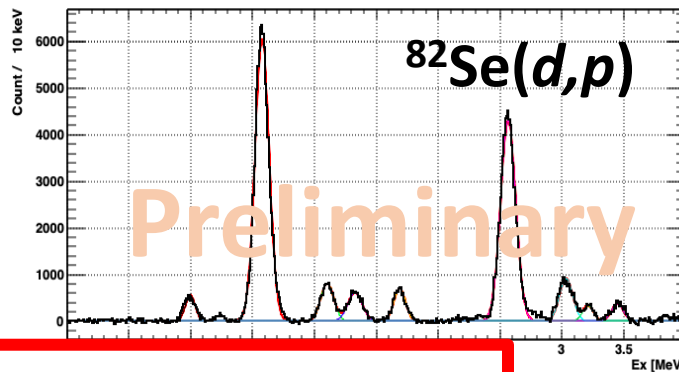
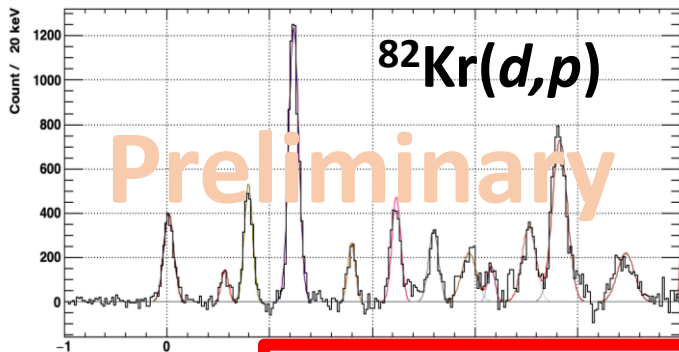
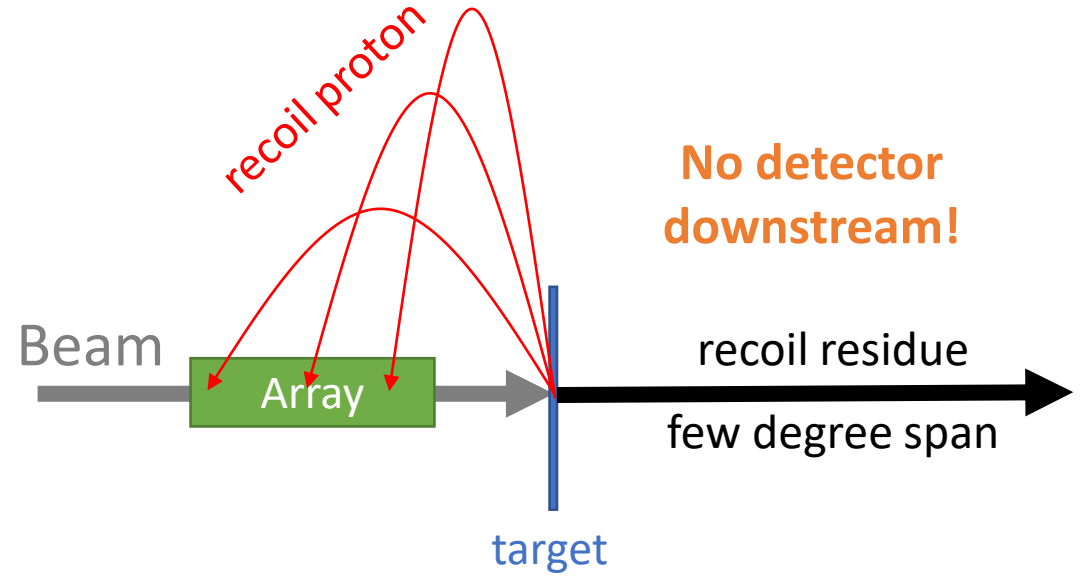
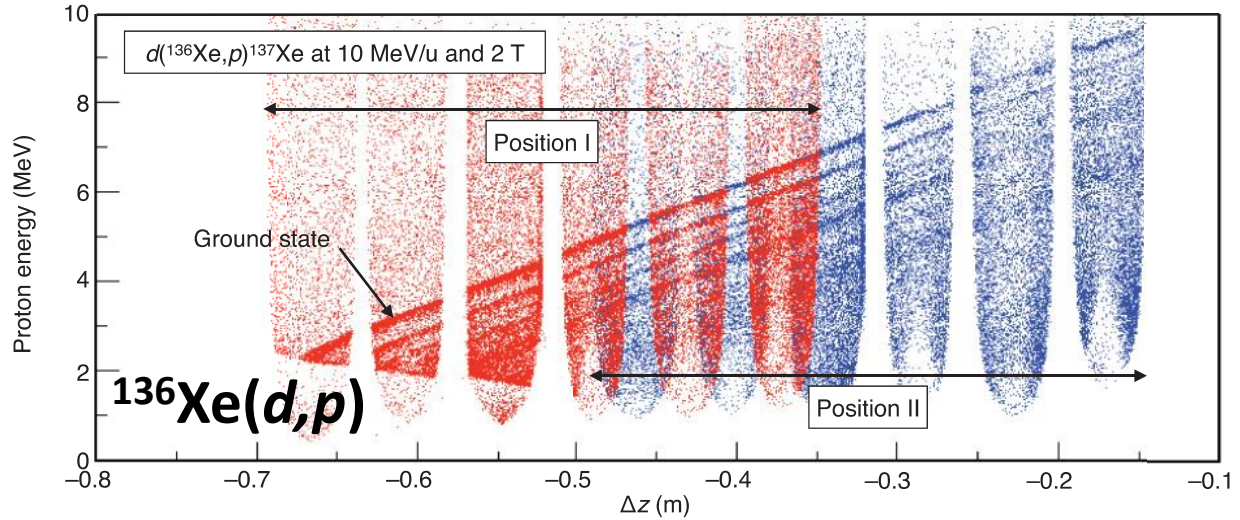


Resolution ~ 100 keV FWHM



Medium Mass Nucleus with HELIOS

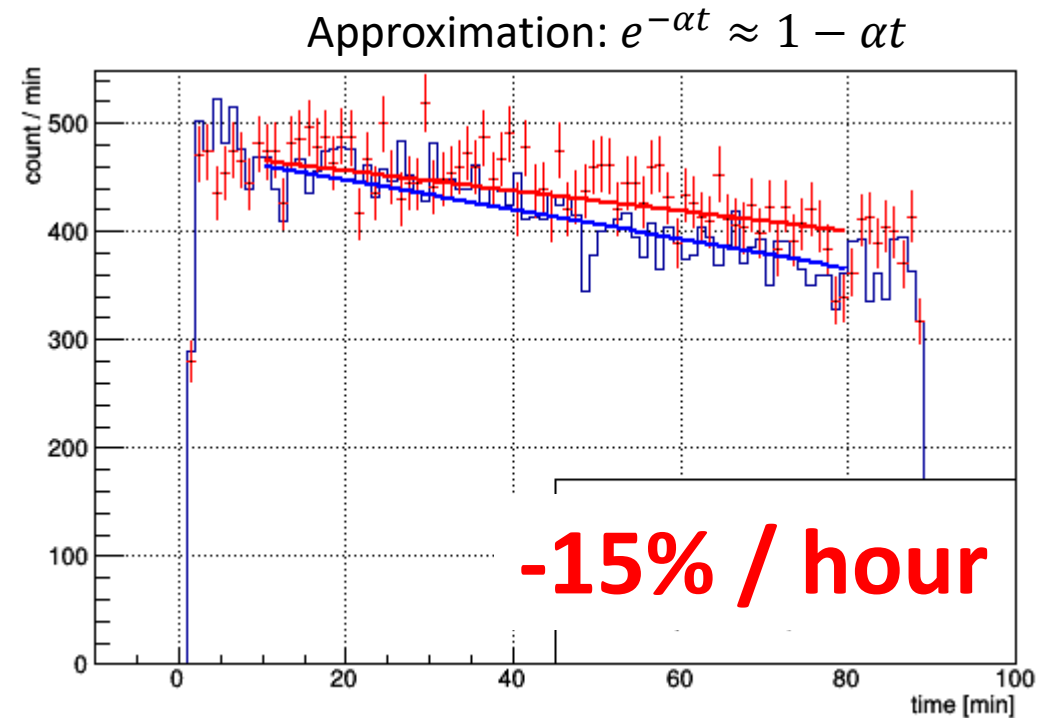
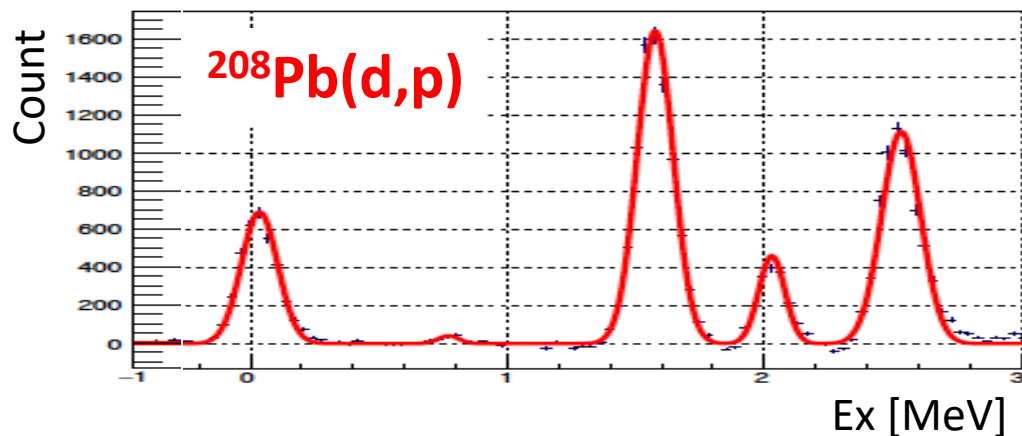
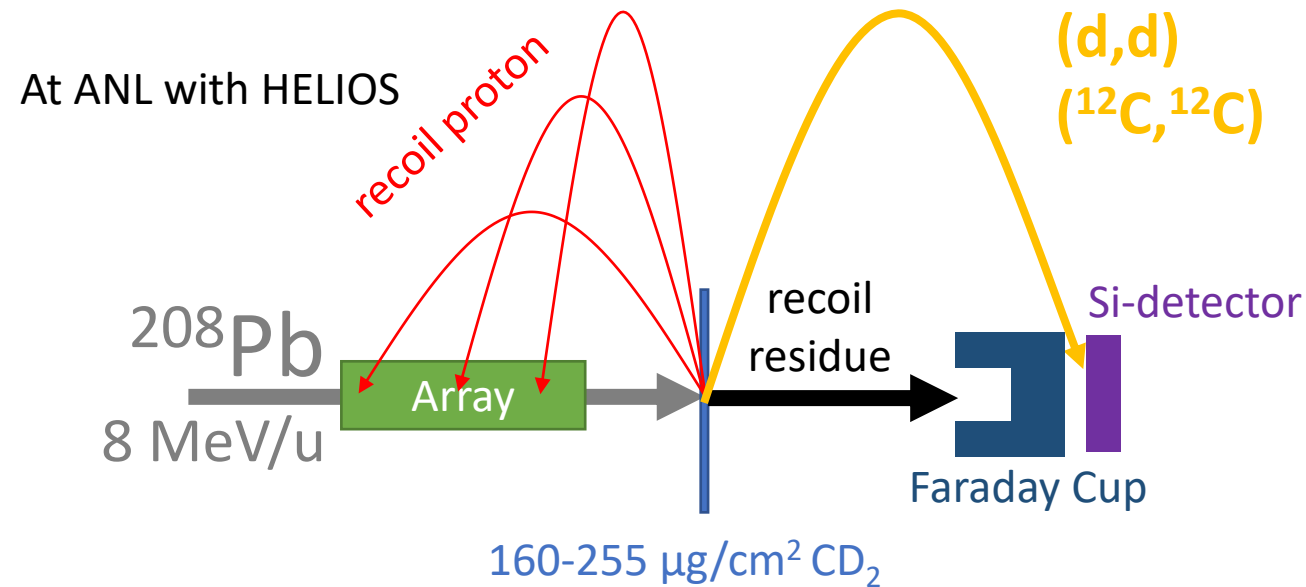
B.P. Kay *et al.*, PRC **84**, 024325 (2011)



Resolution ~ 130 keV FWHM

Medium mass nuclei are OK!!!
 Heavy nuclei should be OK too!!

Target degradation with heavy beam

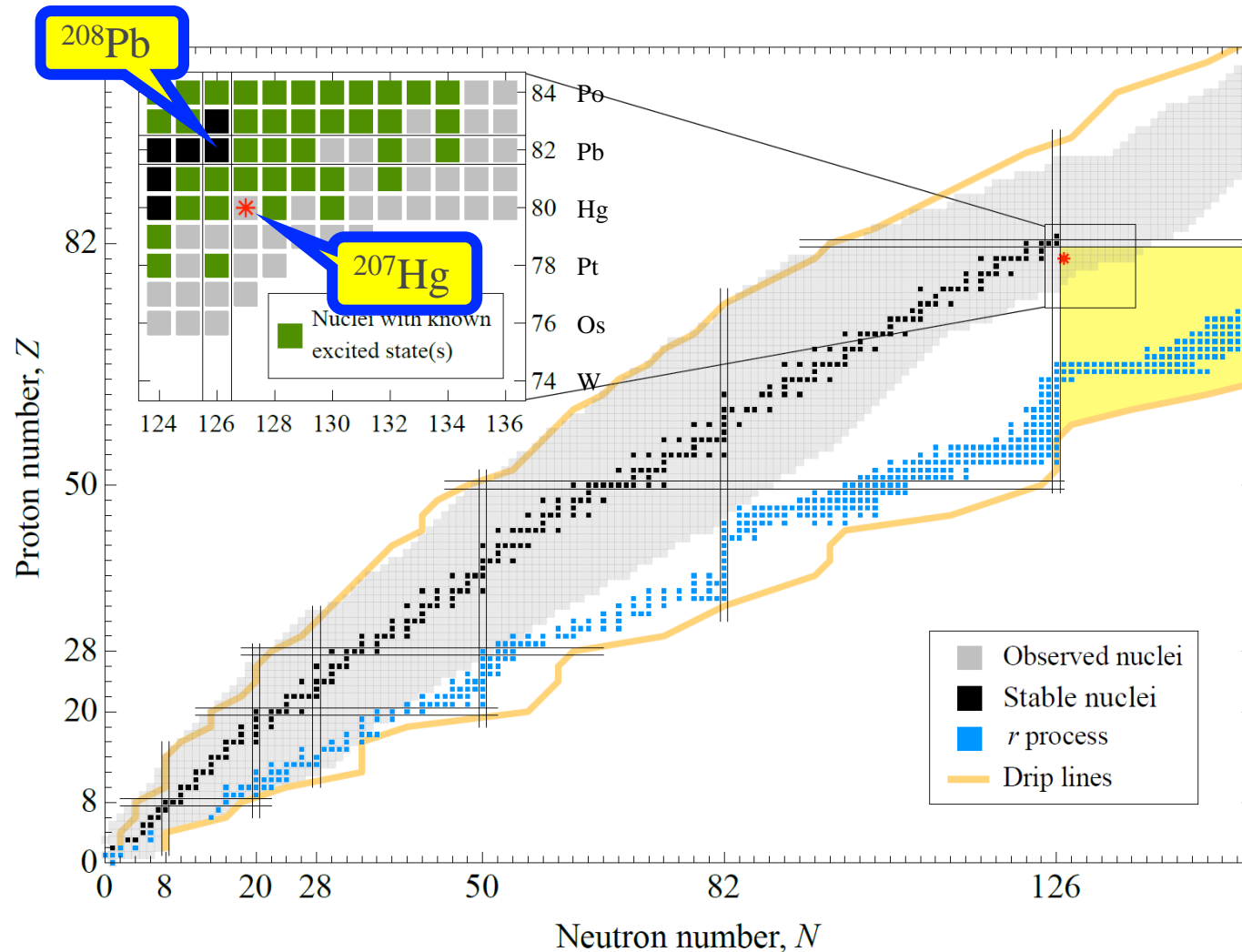


167 $\mu\text{g}/\text{cm}^2$

Beam rate : 8.37×10^6 pps

Total particle: 6.0×10^{10}

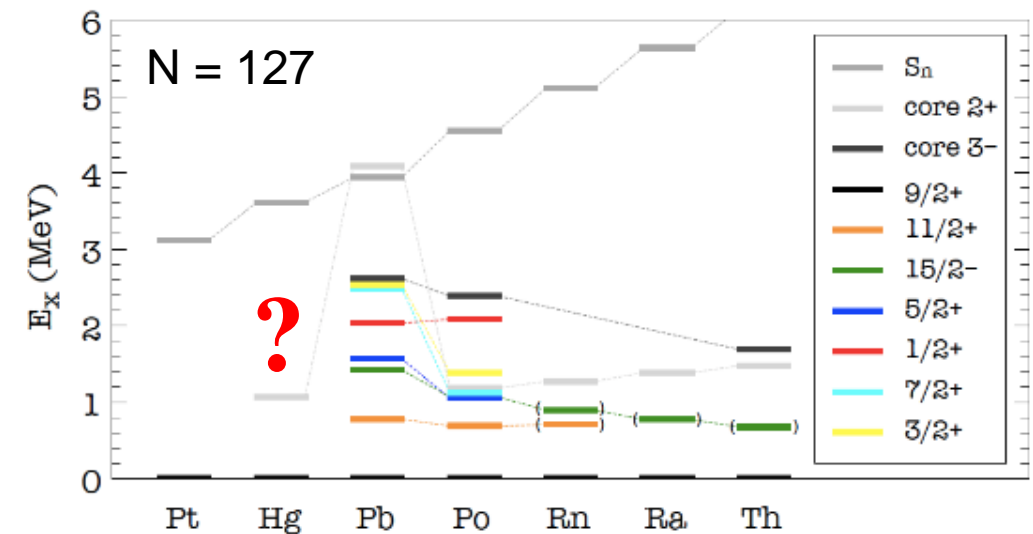
Into Terra Incognita.....



^{208}Pb is a cornerstone to our understanding of the **single-particle structure** of heavy nuclei.

Study of ^{207}Hg

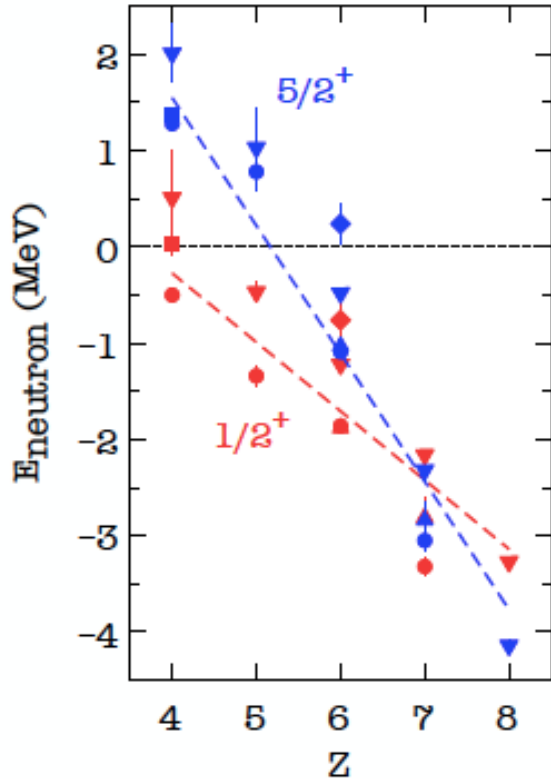
- Toward **r-process** for heavy elements
 - Evolution of shell-structure along $N=127$
 - NO spectroscopy study of isotones below Pb
- **First measurement** on the single-particle state of ^{207}Hg !!!



Weak binding and nuclear structure

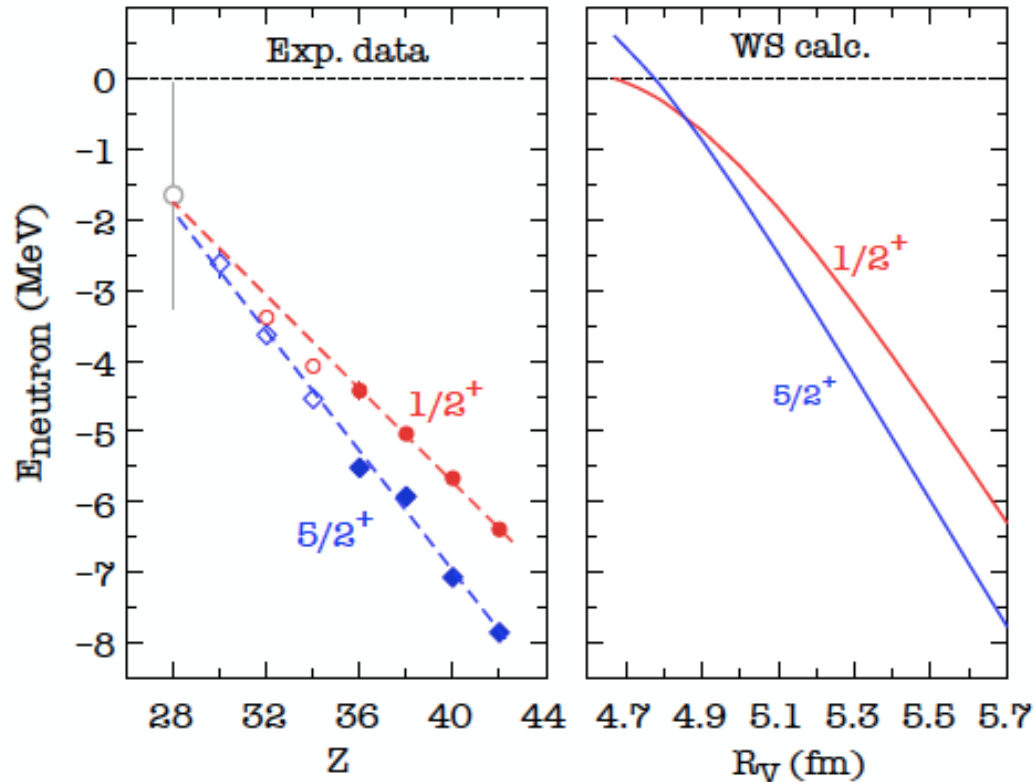
$A \sim 7 - 15$

$1s_{1/2}$ w.r.t $0d_{5/2}$



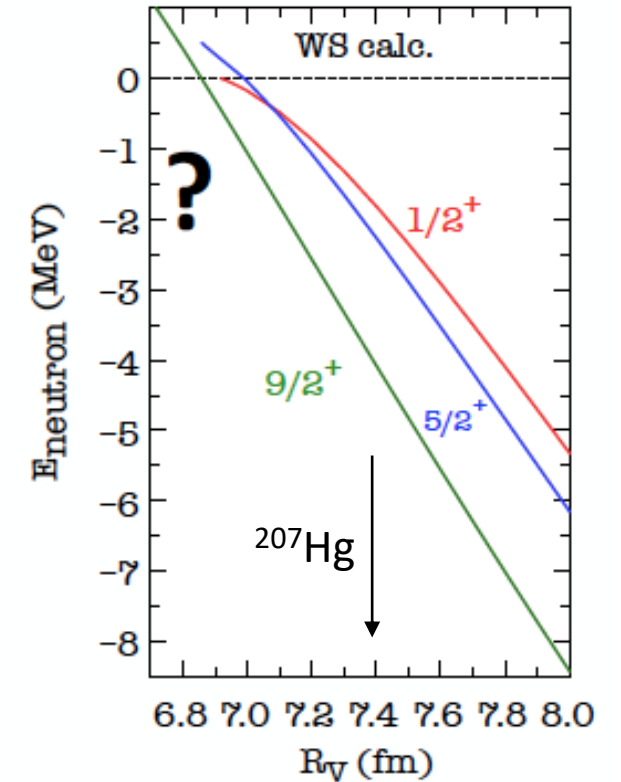
$A \sim 90$

$2s_{1/2}$ w.r.t $1d_{5/2}$



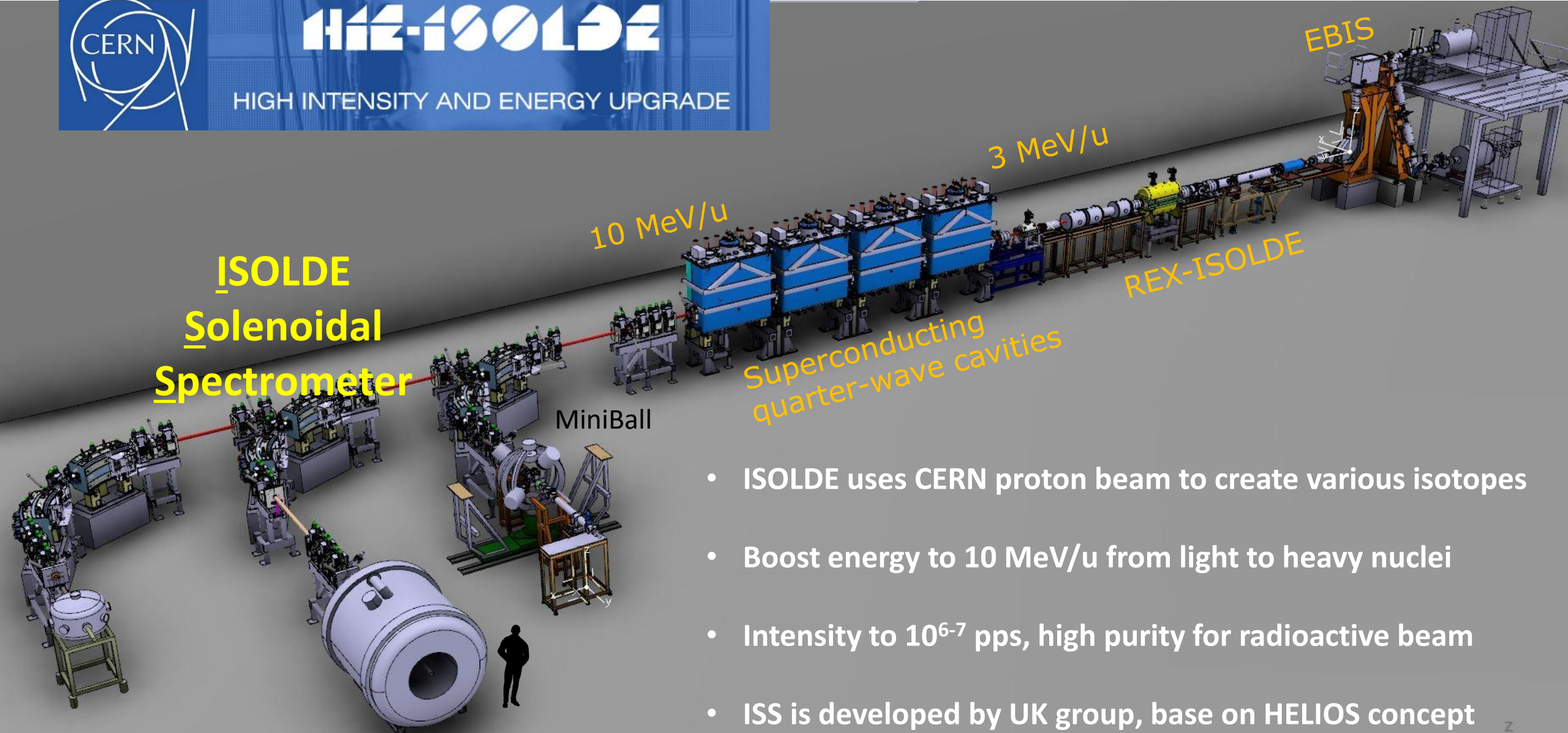
$A \sim 200$

$3s_{1/2}$ w.r.t $2d_{5/2} / 1g_{9/2}$



C. R. Hoffman et al., PRC89(2014)061305

Is it a universal behavior?



ISOLDE Solenoidal Spectrometer

10 MeV/u

3 MeV/u

EBIS

REX-ISOLDE

Superconducting
quarter-wave cavities

MiniBall

- ISOLDE uses CERN proton beam to create various isotopes
- Boost energy to 10 MeV/u from light to heavy nuclei
- Intensity to 10^{6-7} pps, high purity for radioactive beam
- ISS is developed by UK group, base on HELIOS concept

ISOLDE Solenoidal Spectrometer

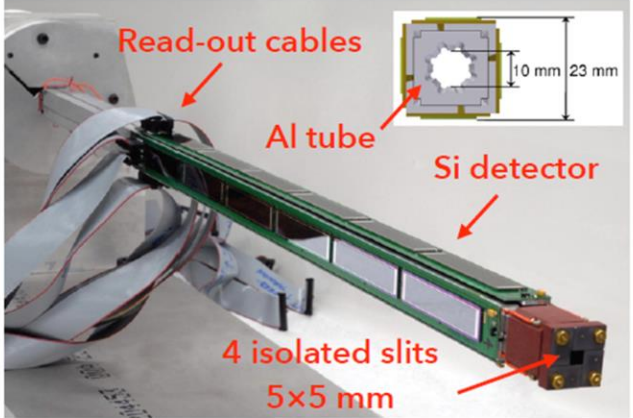
First commission on Sept 2018




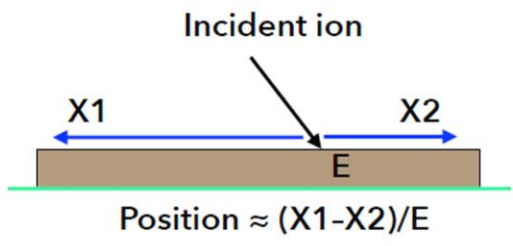
Setup at the time of commission

Both the DAQ and the detector are borrowed from ANL.

heavy beam is used,
radiation damage is large
→ target degradation,
→ A dedicated **8 × 8** type-writer target ladder is used.



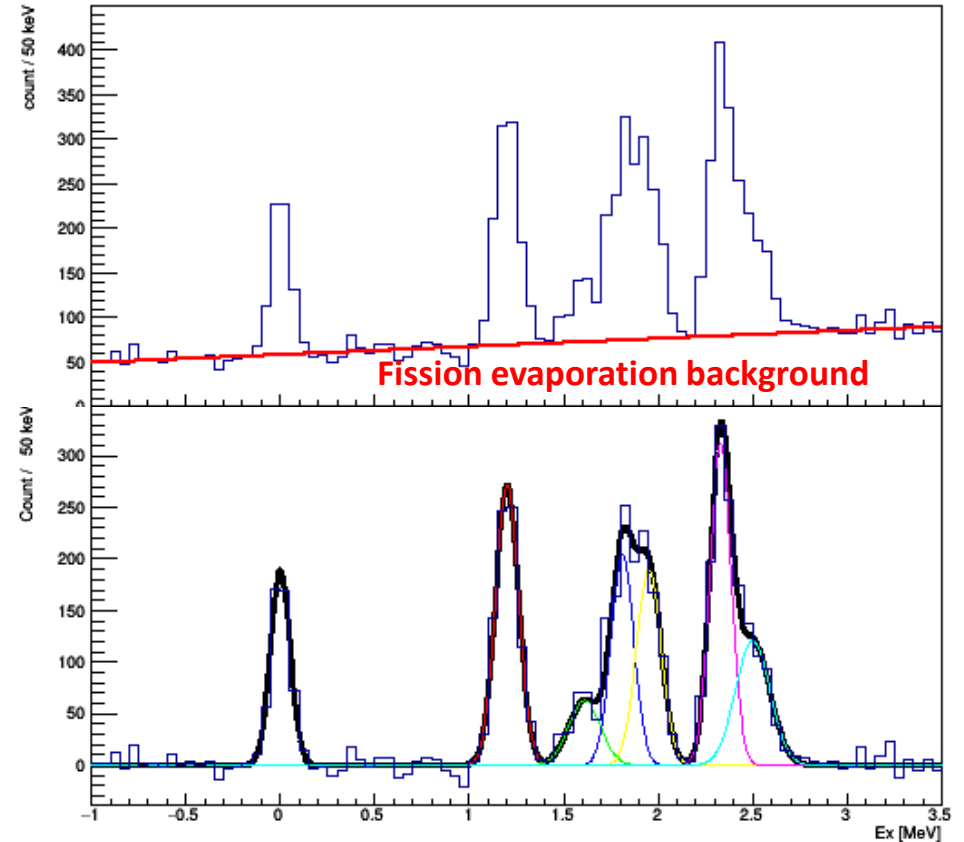
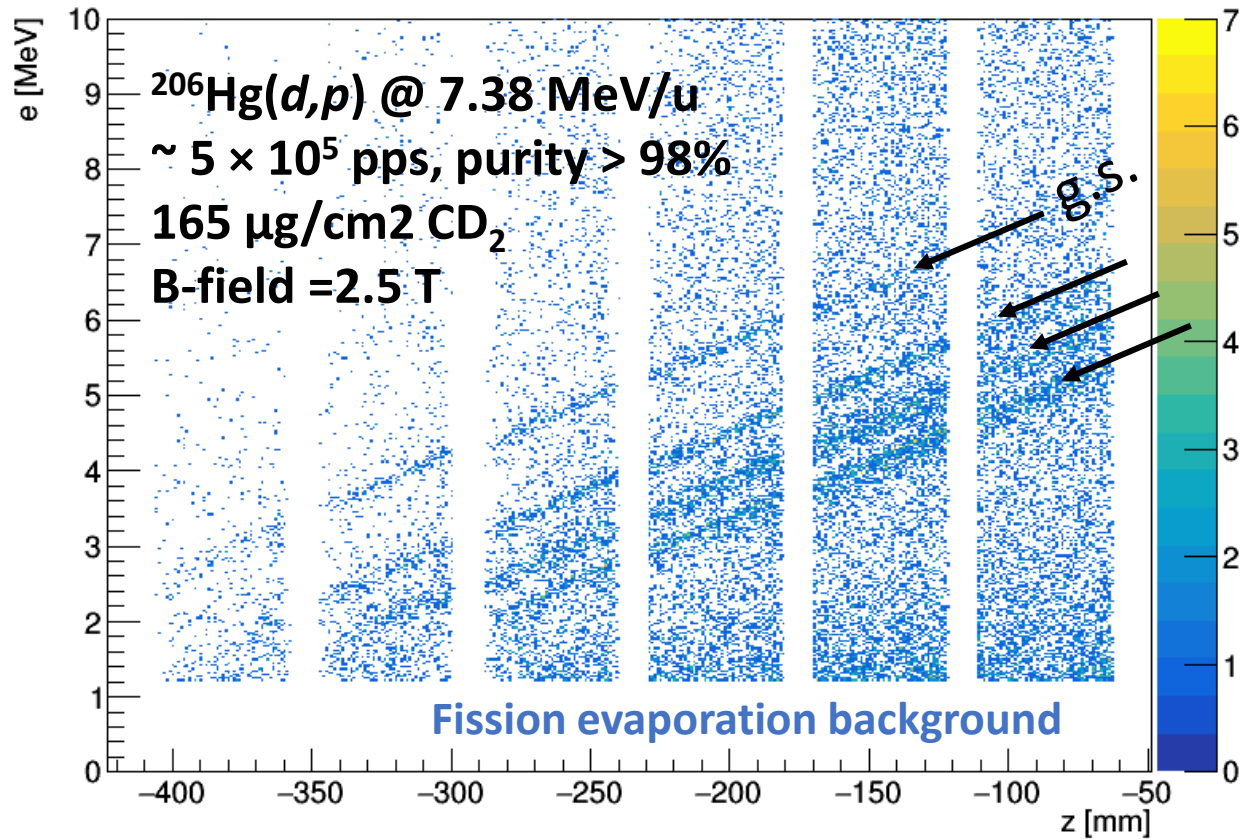
- 4 sides, 6 detectors long
- Detector size, 9×50 mm
- 700- μ m thick (e.g. ~10 MeV protons)
- Φ coverage, **0.48 of 2π**
- $\Omega_{\text{detector}} = 21 \text{ msr}$
- $\Omega_{\text{array}} = 493 \text{ msr}$


Target ladder Mechanical: Russell A. Knaack
Targets: Matthew D. Gott

*J. C. Lighthall et al., Nucl. Instrum. Methods Phys. A **662**, 97 (2010)*

Experimental Result



Resol. ~ 140 keV (FWHM)

Angular Distributions & Spectroscopic Factors

DWBA calculation: Ptolemy

d-channel : A. J. Koning and J. P. Delaroche, Nucl. Phys. A**713**, 231 (2003)

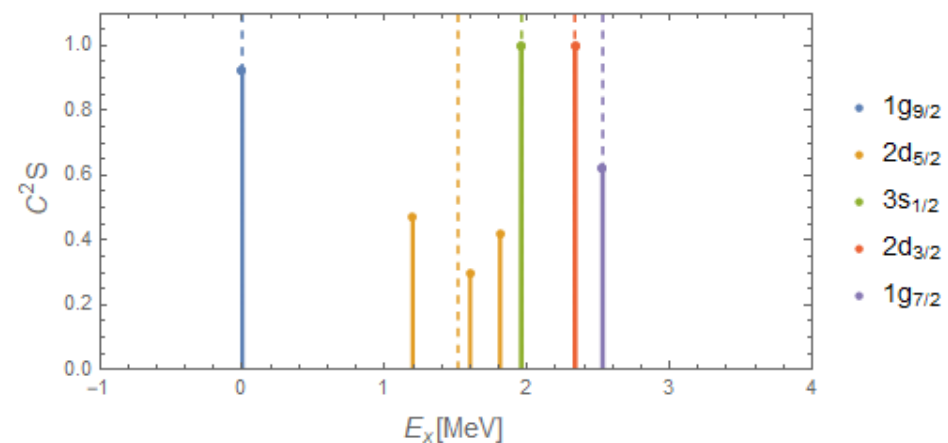
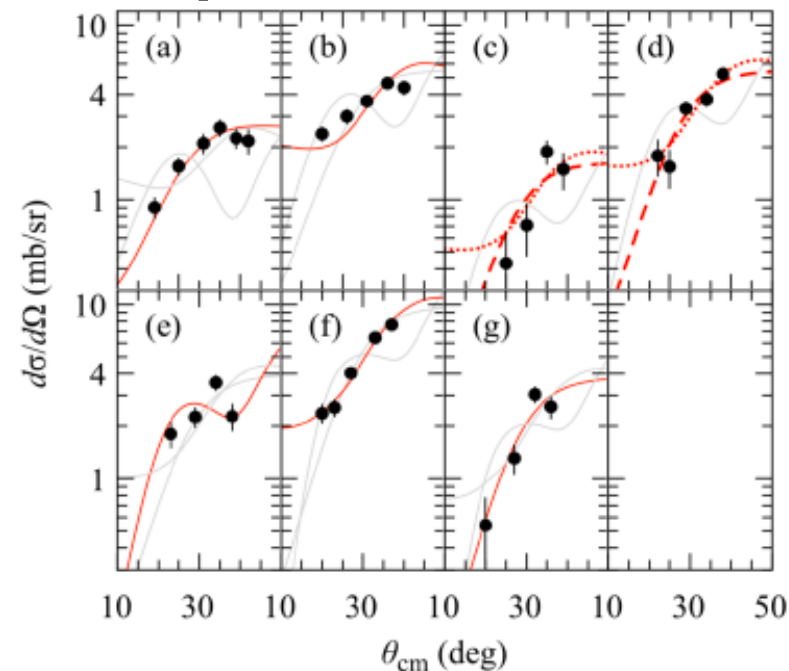
p-channel : H. An and C. Cai, PRC **73**, 054605 (2006)

Bound state: Woods-Saxon, $r_0 = 1.28$ fm, $a_0 = 0.65$ fm

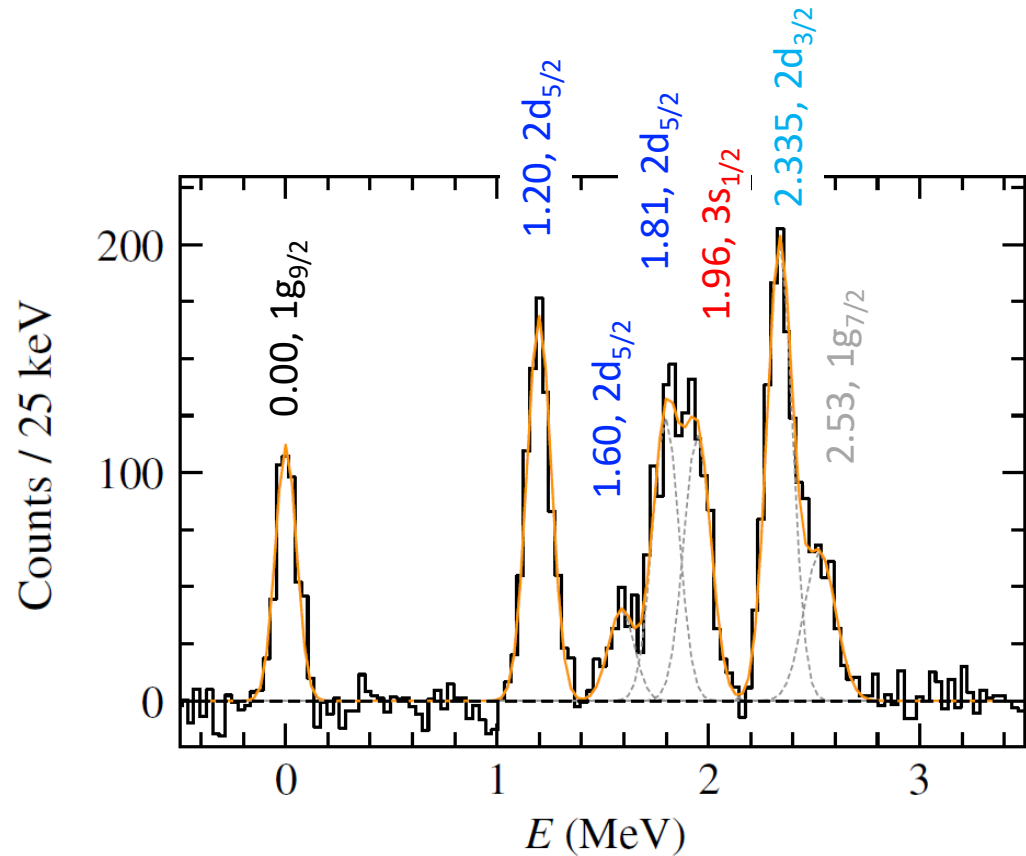
Spin-parity assignment and Spectroscopic factors

E (keV)	ℓ	j^π	nls	S	χ^2/dof
0	4	$9/2^+$	$1g_{9/2}$	0.82(5)	1.3(8)
1197(5)	2	$5/2^+$	$2d_{5/2}$	0.47(6)	2.9(1.1)
1600(45)	4	$9/2^+$	$1g_{9/2}$	0.30(4)	1.5(2)
	2	$5/2^+$	$2d_{5/2}$	0.13(1)	1.4(3)
1810(20)	4	$9/2^+$	$1g_{9/2}$	0.93(12)	1.1(1)
	2	$5/2^+$	$2d_{5/2}$	0.42(3)	1.3(3)
1960(30)	0	$1/2^+$	$3s_{1/2}$	1.00(13)	4.4(2.7)
2335(6)	2	$3/2^+$	$2d_{3/2}$	1.00(7)	1.1(9)
2530(20)	4	$7/2^+$	$1g_{7/2}$	0.62(6)	1.4(2)

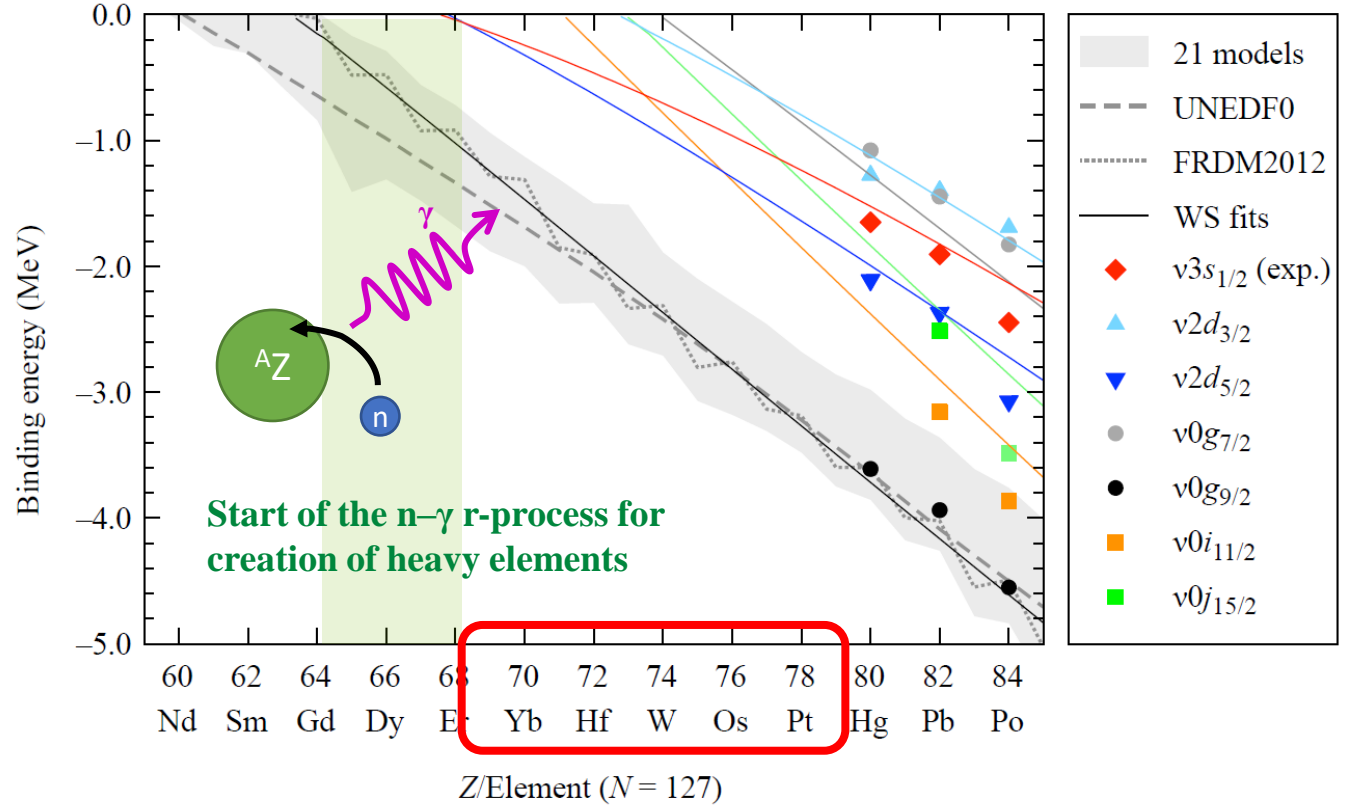
* Normalized to $3s_{1/2}$ state.



Connection to r-process



Woods-Saxon calculations fitted to experimental binding energies of the neutron orbitals at $N = 127$ (^{207}Hg , ^{209}Pb , and ^{211}Po).



Those isotope could be produced in FRIB era.

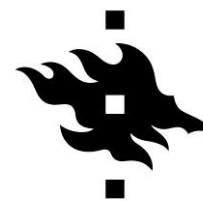
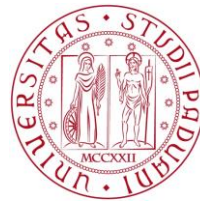
PHYSICAL REVIEW LETTERS **124**, 062502 (2020)

First Exploration of Neutron Shell Structure below Lead and beyond $N = 126$

T. L. Tang,¹ B. P. Kay^{1,*}, C. R. Hoffman,¹ J. P. Schiffer,¹ D. K. Sharp,² L. P. Gaffney,³ S. J. Freeman,² M. R. Mumpower,^{4,5}
 A. Arokiaraj,⁶ E. F. Baader,³ P. A. Butler,⁷ W. N. Catford,⁸ G. de Angelis,⁹ F. Flavigny,^{10,11} M. D. Gott,¹ E. T. Gregor,⁹
 J. Konki,³ M. Labiche,¹² I. H. Lazarus,¹² P. T. MacGregor,² I. Martel,⁷ R. D. Page,⁷ Zs. Podolyák,⁸ O. Poleshchuk,⁶
 R. Raabe,⁶ F. Recchia,^{13,14} J. F. Smith,¹⁵ S. V. Szwece,^{16,17} and J. Yang⁶



The University of Manchester



UNIVERSITÀ
DEGLI STUDI
DI PADOVA



UNIVERSITY OF JYVÄSKYLÄ



Outlook

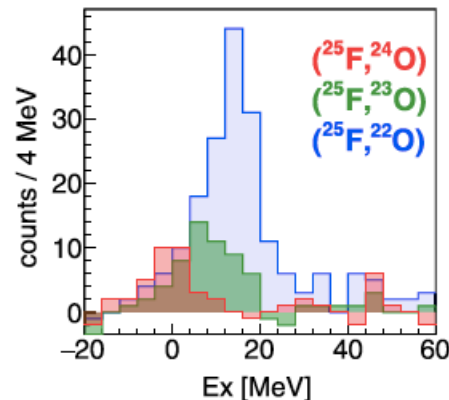
The future of :

- Single-particle spectroscopy
- Solenoidal spectrometers

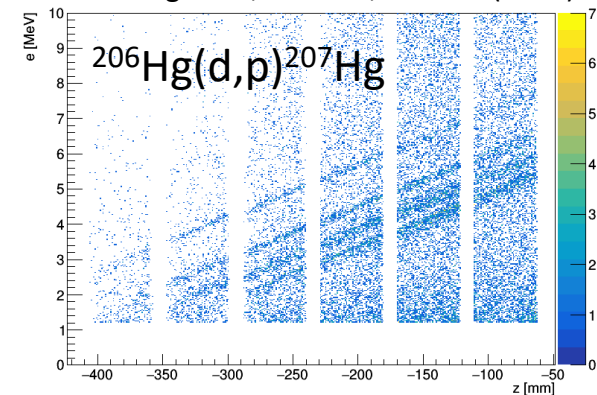
More Applications

There are new *tricks* for single-particle spectroscopy.

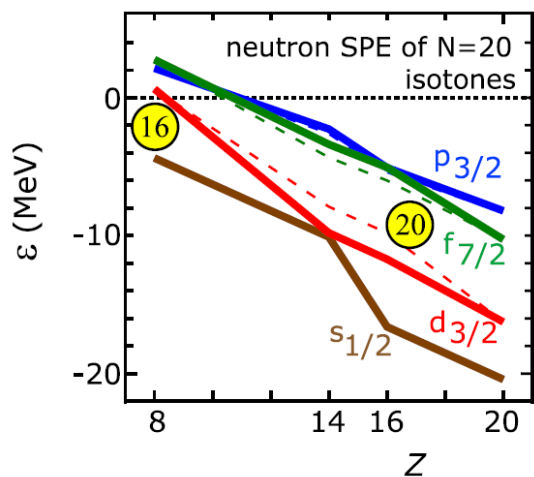
T.L.Tang et al., PRL **124**, 212502 (2020)



T.L.Tang et al., PRL **124**, 062502 (2020)

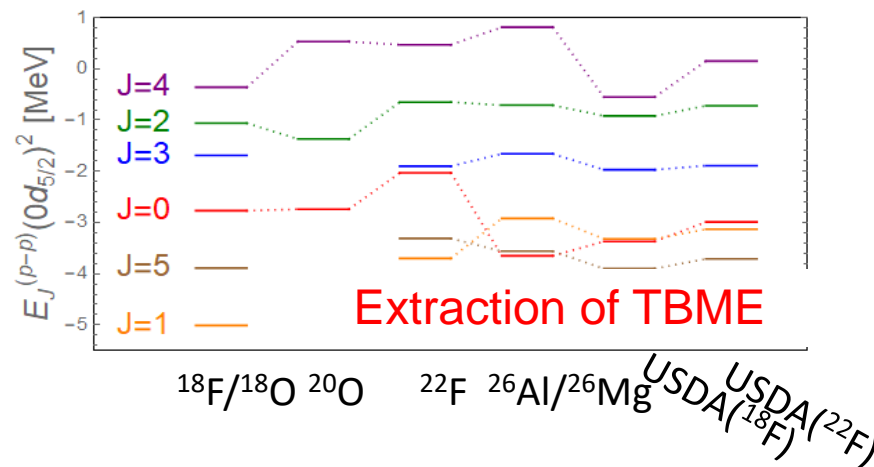


T. Otsuka, Phys. Scr. **T152** (2013) 014007

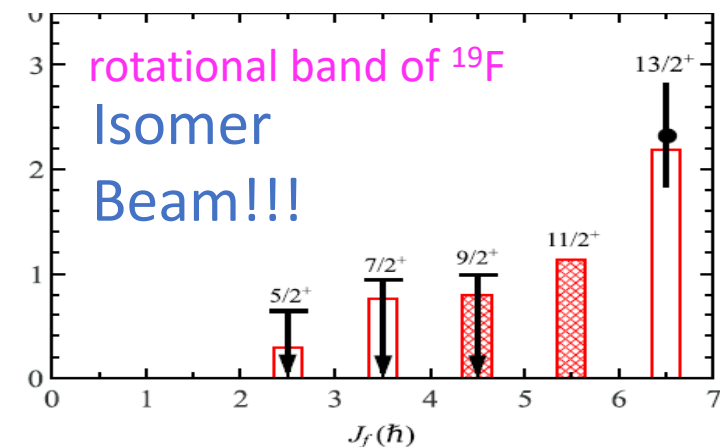


Study of ESPE and shell evolution

J. Chen et al., PRC **98**, 014325 (2018)



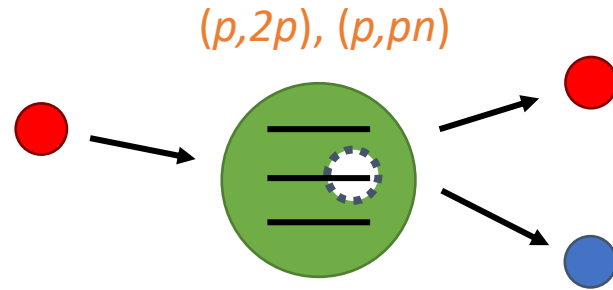
D. Santiago-Gonzalez et al., PRL **120**, 122503 (2018)



Connection between single-particle and collective motion.

Above are only few applications, a lot MORE !!!

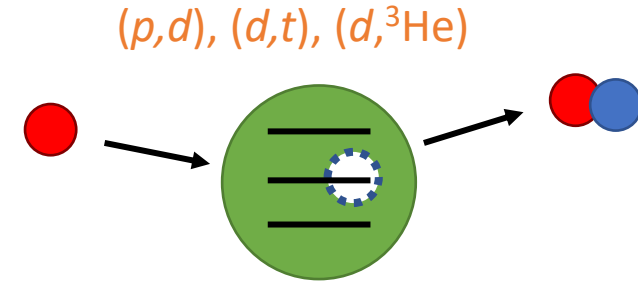
The needs of *Fast* and *Slow* beams reactions



Fast beam knockout

- Deep inelastic scattering.
- No Q-value limitation.
- Simpler reaction mechanism (Impulse approximation)
- High energy \rightarrow less energy loss
 - Solid polarized proton target
 - Thick target technique
 - Beam event-by-event PID & tracking
 - Not everything in vacuum

Study the **fullness** of orbital

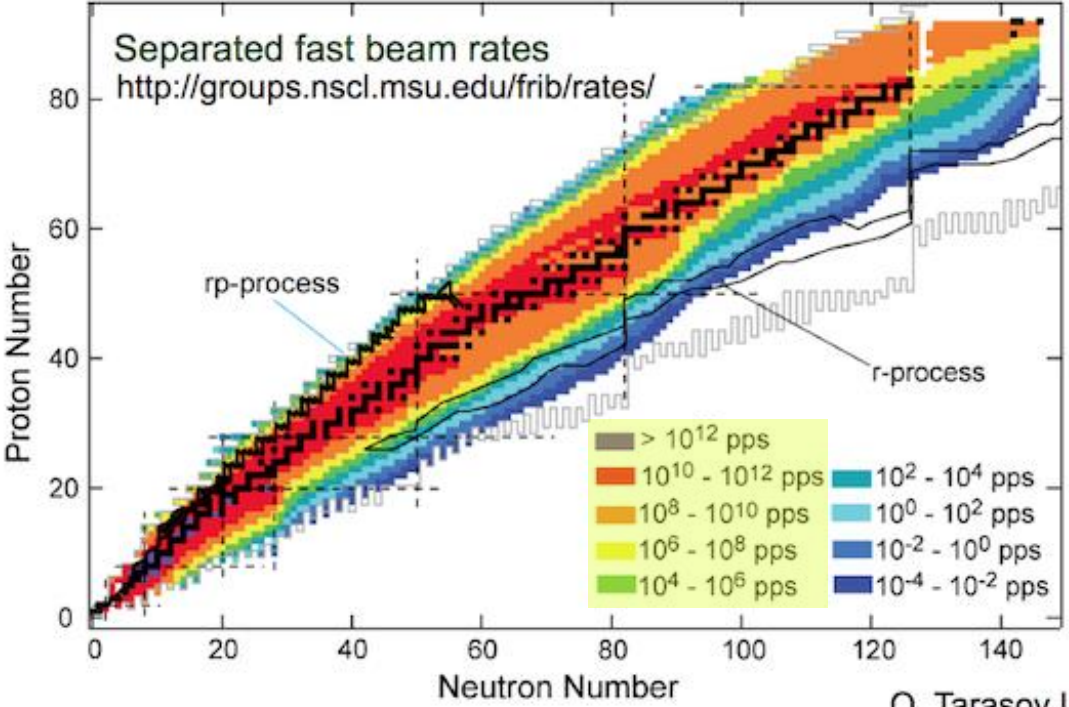


Slow beam removal

- On nuclear surface
- Sometimes limited by Q-value.
 - \rightarrow Beam energy $>$ - Q-value
- Easy for neutron removal
- Reaction mechanic is well-understood

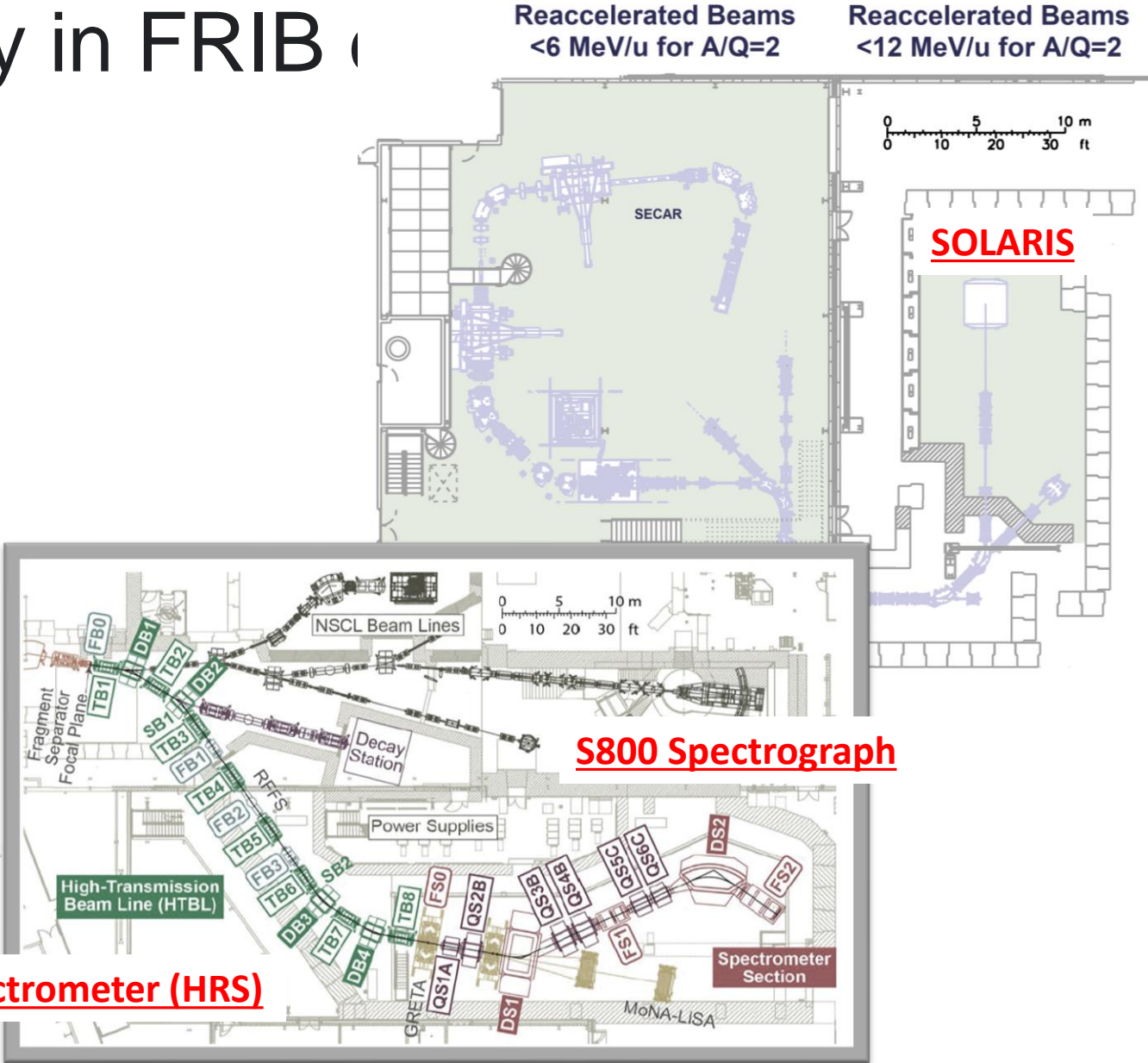
Both can extend to *Quasi-particle spectroscopy!!!*

Single-particle spectroscopy in FRIB



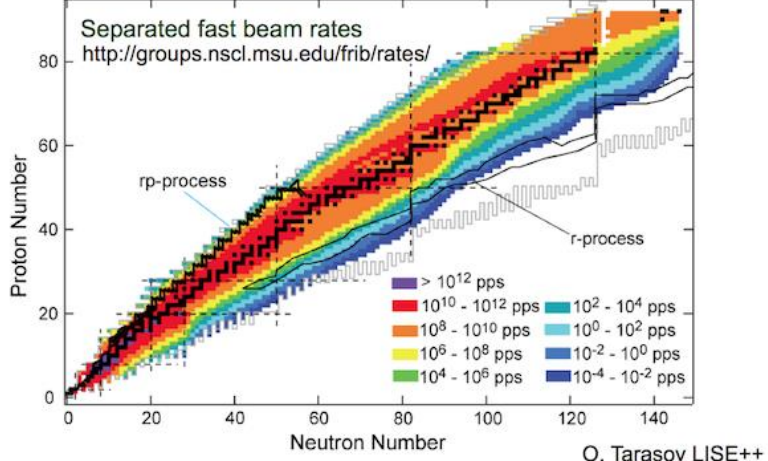
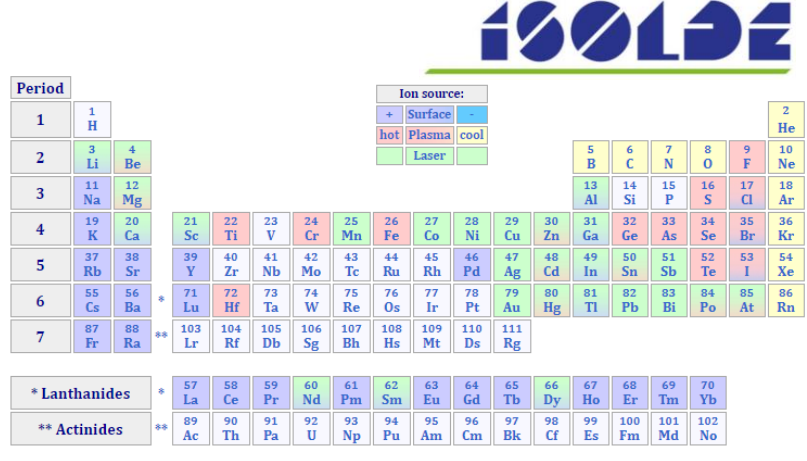
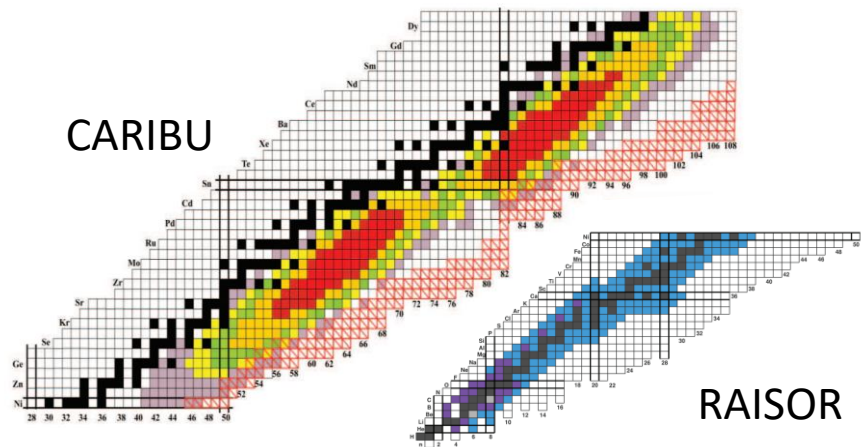
O. Tarasov LISE++

High Rigidity Spectrometer (HRS)

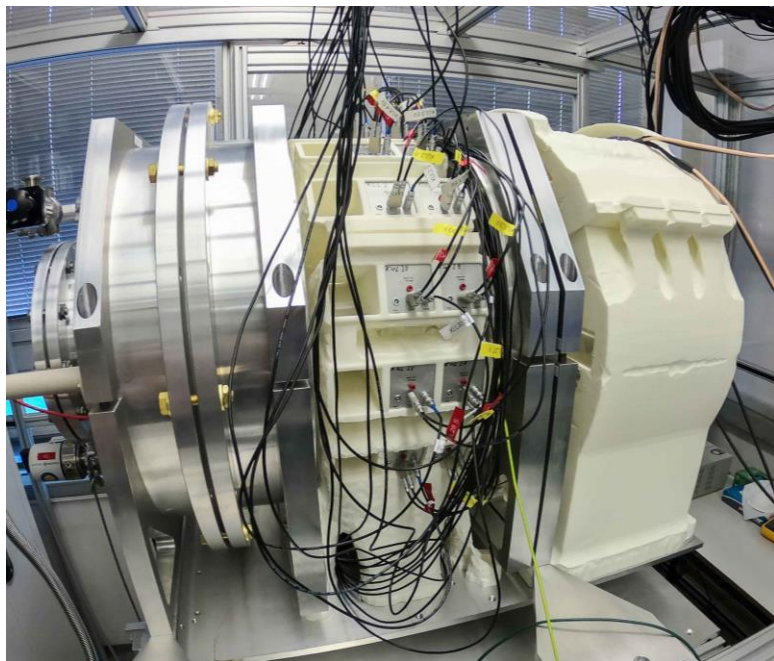


S800 Spectrograph

Future of the Solenoidal Spectrometers



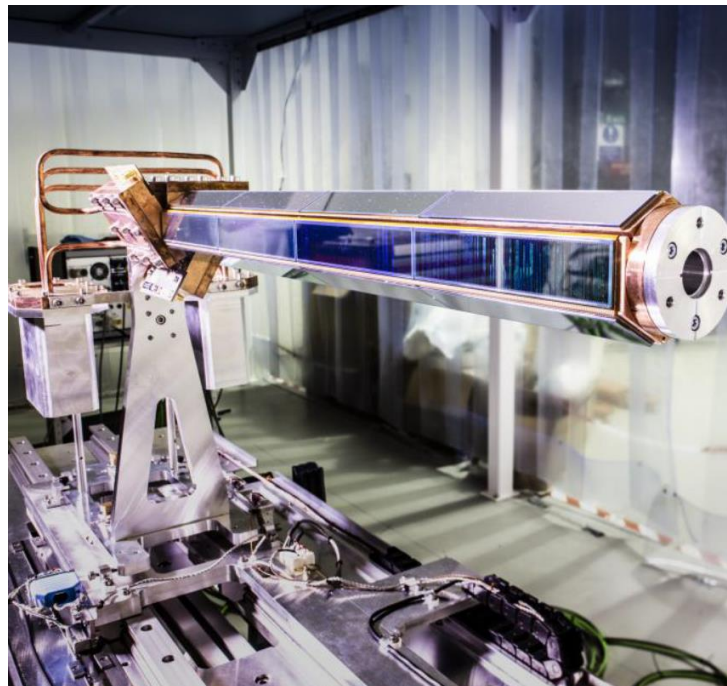
New Devices for Solenoidal Spectrometers



SpecMat : Charged particle
+ Gamma-ray detector
Lead by **Oleksii Poleshchuk**, KU Leuven

X-Y Sensitive Si-detector

Lead by **Robert Page**,
University of Liverpool



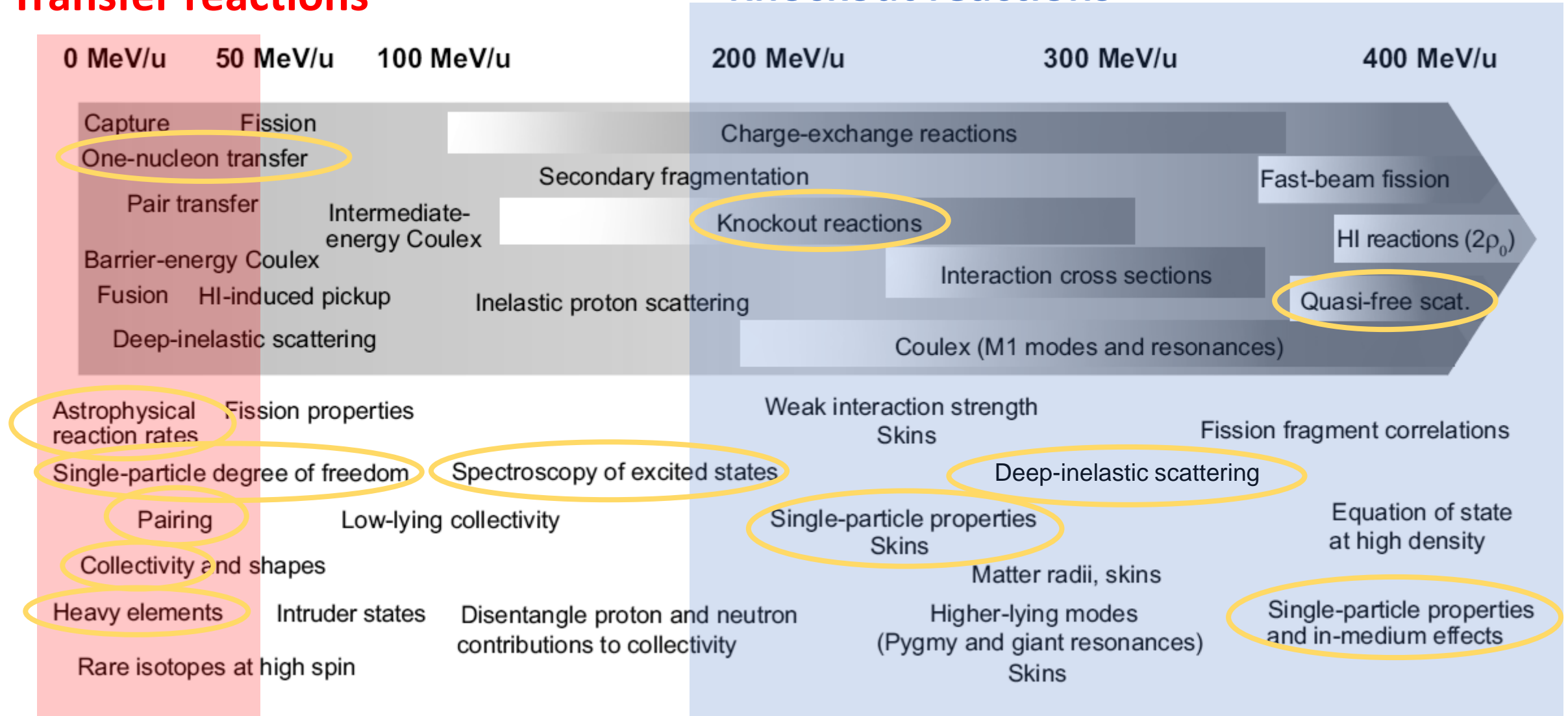
SOLARIS

AT-TPC, Lead by **Daniel Bazin**, NSCL



Transfer reactions

Knockout reactions



Taken from : FRIB400 – The scientific case for the 400 MeV/u Energy Upgrade of FRIB

Summary

Single-particle spectroscopy is an active tool with many applications

- most intuitive picture of how nucleon assembly.

But still many new frontiers, 2 examples:

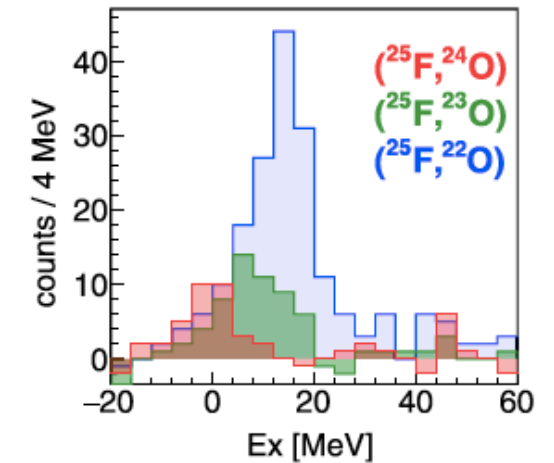
- **Fast beam** ($p,2p$) : study neutron shell of ^{25}F
- **Slow beam** (d,p) : single-neutron spectroscopy of ^{207}Hg

There are more applications!

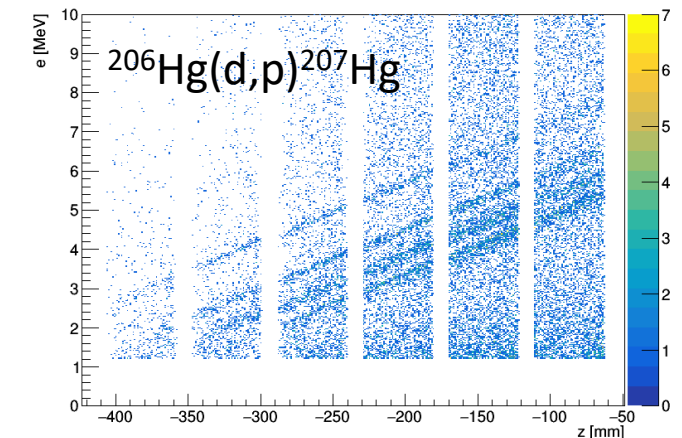
Solenoidal spectrometers are around the world!



T.L.Tang et al., PRL **124**, 212502 (2020)



T.L.Tang et al., PRL **124**, 062502 (2020)



Thank you for your attention!!!